

**Длинные нелинейные стоячие волны
в протяженном бассейне с пологими берегами:
теория и эксперимент**

С.Ю.Доброхотов, В.А.Калиниченко, Д.С.Миненков, В.Е.Назайкинский

Институт проблем механики им.А.Ю.Ишлинского РАН
и Московский физико-технический институт

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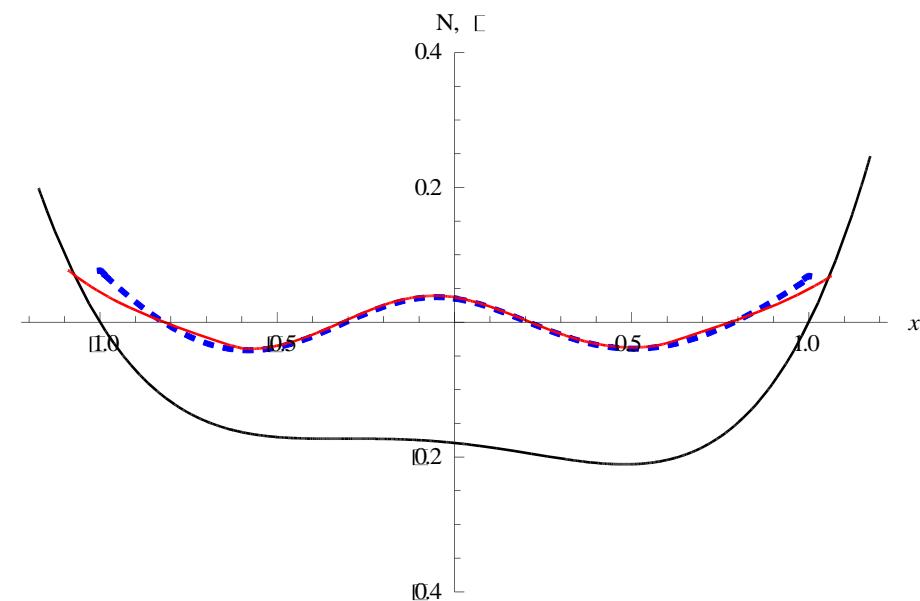
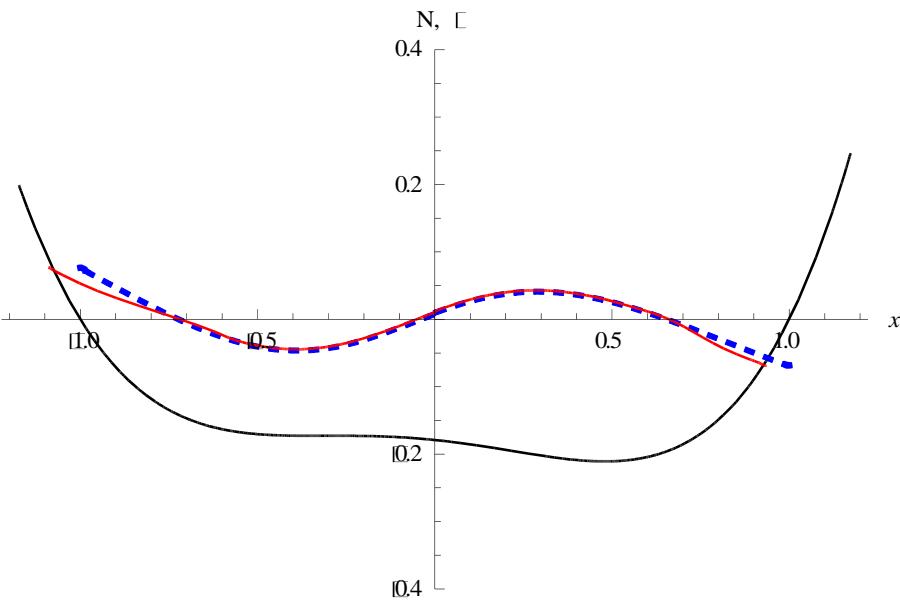
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Shallow water: two shores

$$\eta_t + ((D(x) + \eta)u)_x = 0, \quad u_t + g\eta_x + uu_x = 0.$$

Depth $D(a) = D(b) = 0$ $D'(a) \neq 0$, $D'(b) \neq 0$

Boundary condition at two variable boundaries $x(t)$: $\eta(x_a(t), t) + x_a(t) = 0$



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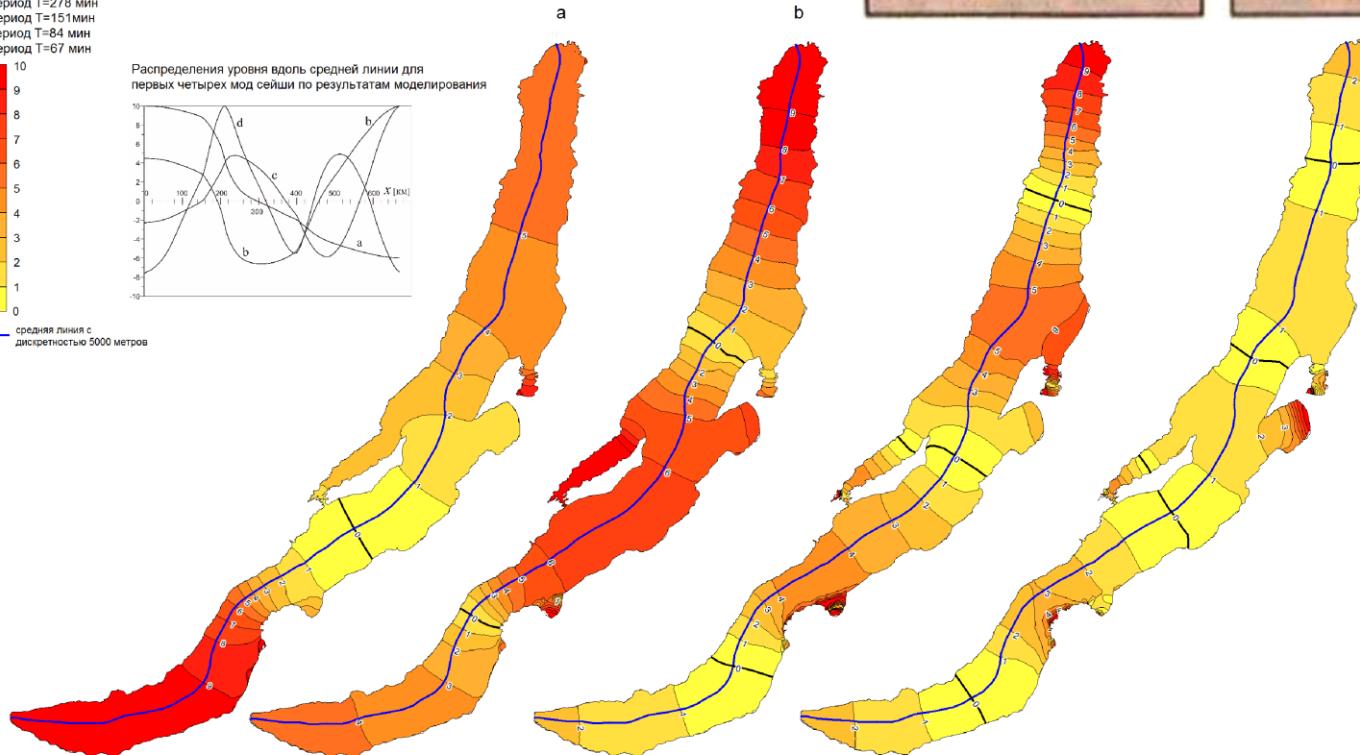
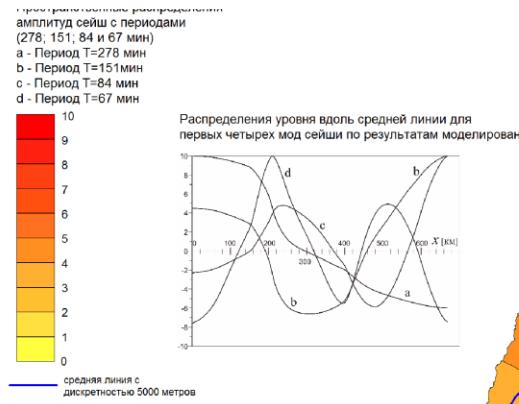
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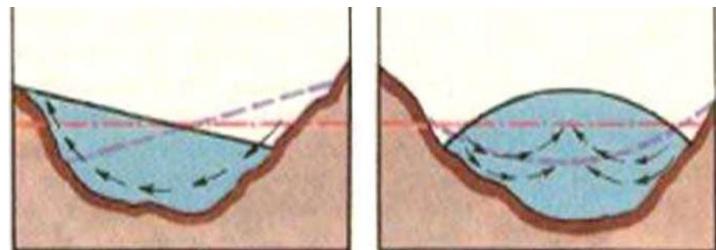
Motivation: Seiches – standing waves

A **seiche** is a standing wave in an enclosed or partially enclosed body of water.

The Baikal Lake seiches



One- and two-nodes seiches



МАСШТАБ 1: 2 500 000

Linearized shallow water: two shores

$$N_\tau + (D(y)U)_y = 0, \quad U_\tau + N_y = 0.$$

Depth $D(a) = D(b) = 0$

No boundary conditions,

$\therefore D'(a) \neq 0, \quad D'(b) \neq 0$

Finite energy condition: $E^2 = \|N\|^2 + \|\sqrt{D(y)} U\|^2 < \infty$

Asymptotics.

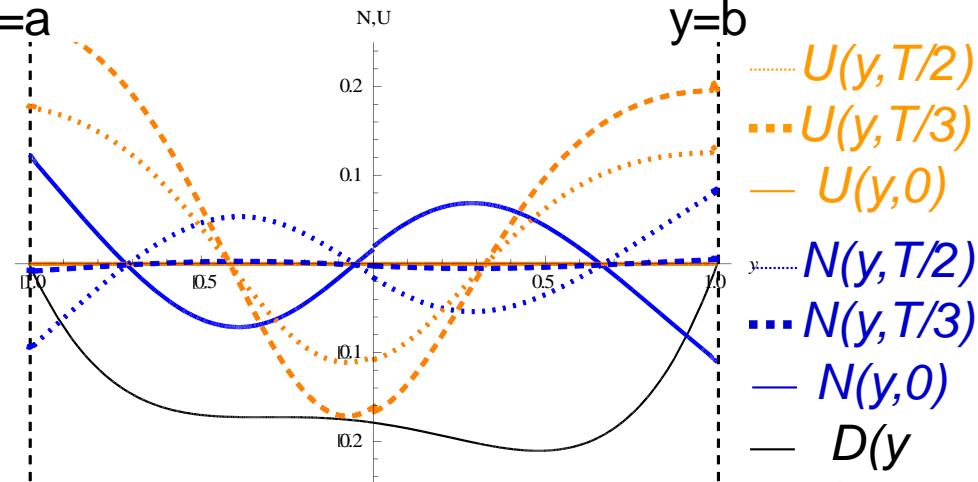
Phase: $S(x, y) = \int_x^y \frac{dy}{\sqrt{D(y)}}, \quad a \leq x \leq y \leq b.$

Quantization: $w_n = \frac{\pi}{S(a, b)} (\frac{1}{2} + n)(1 + O(n^{-1})), \quad n \in \mathbb{N}.$

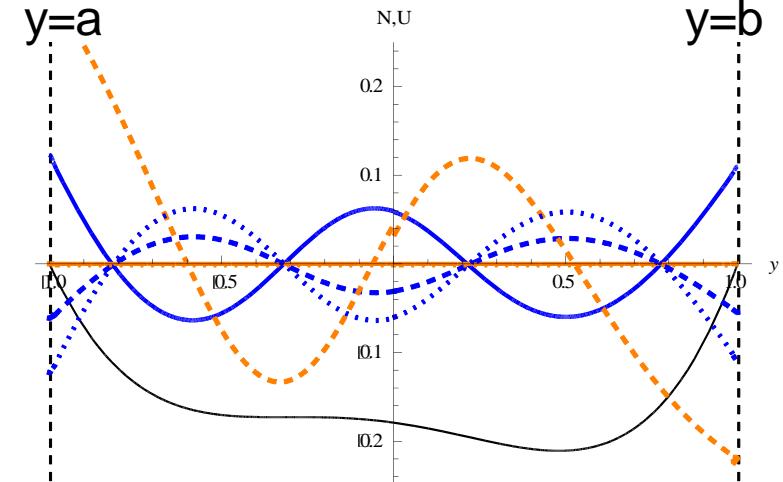
$$N_a(y, \tau) = c \cos(w_n \tau + \tau_0) \mathbf{J}_0(w_n S(a, y)) \left(\frac{S(a, y)}{\sqrt{D(y)}} \right)^{1/2} \quad y \in [a, b - \delta]$$

$$N_b(y, \tau) = c(-1)^n \cos(w_n \tau + \tau_0) \mathbf{J}_0(w_n S(y, b)) \left(\frac{S(y, b)}{\sqrt{D(y)}} \right)^{1/2} \quad y \in [a + \delta, b]$$

$y=a$



$y=a$



Linearized shallow water: shore and vertical wall

$$N_\tau + (D(y)U)_y = 0, \quad U_\tau + N_y = 0 \quad \text{Depth} \quad D(a) = 0, D'(a) \neq 0$$

Boundary condition at

$y=b$:
Finite energy condition:

$$U(b, \tau) = 0$$

$$E^2 = \|N\|^2 + \|\sqrt{D(y)} U\|^2 < \infty$$

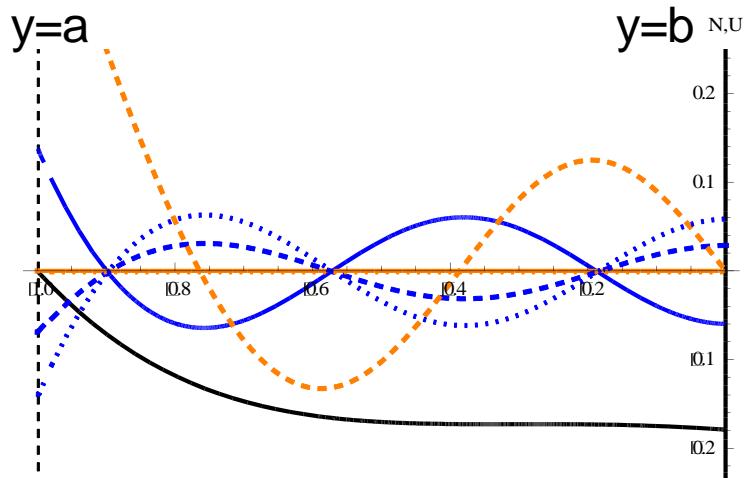
Asymptotics.

$$\text{Phase: } S(x, y) = \int_x^y \frac{dy}{\sqrt{D(y)}}, \quad a \leq x \leq y \leq b.$$

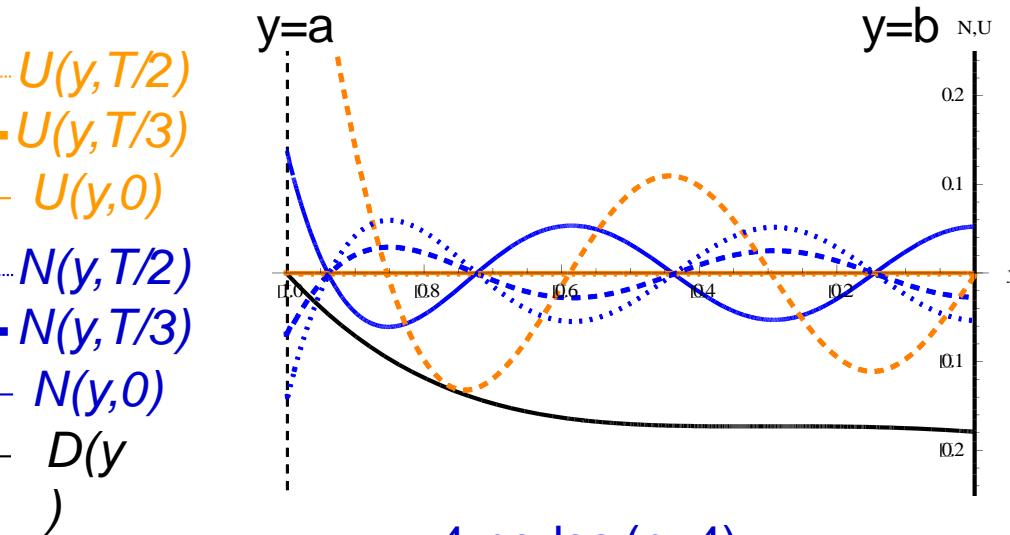
$$\text{Quantization: } w_n = \frac{\mu_n}{S(a, b)} \approx \frac{\pi}{S(a, b)} \left(\frac{1}{4} + n \right) \quad \mathbf{J}_1(\mu_n) = 0, \quad n \in \mathbb{N}$$

$$N_a(y, \tau) = c \cos(w_n \tau + \tau_0) \mathbf{J}_0(w_n S(a, y)) \left(\frac{S(a, y)}{\sqrt{D(y)}} \right)^{1/2} (1 + O(w_n^{-3/2})),$$

$$U_a(y, \tau) = c \sin(w_n \tau + \tau_0) \mathbf{J}_1(w_n S(a, y)) \left(\frac{S(a, y)}{\sqrt{D(y)}} \right)^{1/2} \frac{1}{\sqrt{D(y)}} (1 + O(w_n^{-3/2}))$$



3-nodes (n=3)



4-nodes (n=4)

**The main defect in the linear model:
it does not describe the splash (run up)**

⇒ Carrier-Greenspan transform or its asymptotic modification near the beach

Shallow water: sloping bottom, Carrier—Greenspan

$$\eta_t + ((D(x) + \eta)u)_x = 0, \quad u_t + g\eta_x + uu_x = 0. \quad \begin{matrix} \text{Depth } D(x) = \gamma(x - a) \\ \vdots \end{matrix}$$

Boundary condition at left variable boundary $x(t)$: $\eta(x_a(t), t) + x_a(t) = 0$

Boundary condition at $x=b$: $u|_{x=b} = 0$

Without loss of generality: $g=1$, $a=0$, $\gamma = 1$

$b=1$,

Carrier—Greenspan transform.

$$t = \tau + U, \quad x = y - N + U^2/2, \quad \eta = N - U^2/2, \quad u = U \iff$$

$$\tau = t - u, \quad y = x + \eta - u^2/2, \quad N = \eta + u^2/2, \quad U = u.$$

$$J \equiv \frac{\partial(\tau, y)}{\partial(t, x)} = 1 + \eta_x - u_t - \eta_x u_t + \eta_t u_x, \quad J^{-1} \equiv \frac{\partial(t, x)}{\partial(\tau, y)} = 1 - N_y + U_\tau - N_y U_\tau + N_\tau U_y + U U_y$$

Theorem (C—G).

If $J > 0$, $J < \infty$

then shallow water

is equivalent to linearized SW:

$$\begin{pmatrix} \frac{\partial v}{\partial t} + \frac{\partial[\eta + v^2/2]}{\partial x} \\ \frac{\partial \eta}{\partial t} + \frac{\partial[(\eta + x)v]}{\partial x} \end{pmatrix} = \frac{\partial(\tau, y)}{\partial(t, x)} \begin{pmatrix} \frac{\partial U}{\partial \tau} + \frac{\partial N}{\partial y} \\ \frac{\partial N}{\partial \tau} + \frac{\partial[yU]}{\partial y} \end{pmatrix} = 0$$

$x_a(t)$ becomes fixed boundary: $y_a(t)=0$. Finite energy condition or smth alike is required.

Boundary condition at $x=b$ becomes nonlin: $U(Y(\tau), \tau) = 0, \quad Y(\tau) = b + N(Y(\tau), \tau)$

Reduced Carrier-Greenspan transformation (simplification):

$$D(y) = D(x) + \eta(x, t)\rho(x), \tau = t, N(y, \tau) = \eta(x, t), U(y, \tau) = u(x, t),$$

here $\rho(x)$ is a cut-off function, $\rho = 1$ near x^0 : $D(x) = 0$ and $\rho = 0$ outside of the neighborhood of the point x^0 .

The main property: **the boundary becomes fixed**

Shallow water: two shores

$$\eta_t + ((D(x) + \eta)u)_x = 0, \quad u_t + g\eta_x + uu_x = 0.$$

Depth $D(a) = D(b) = 0$ $D'(a) \neq 0$, $D'(b) \neq 0$

Boundary condition at two variable boundaries $x(t)$: $\eta(x_a(t), t) + x_a(t) = 0$

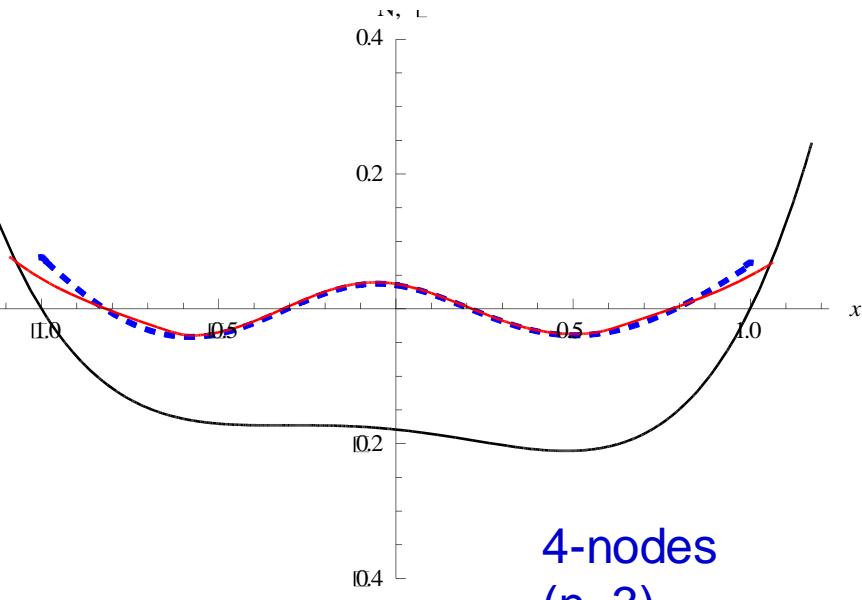
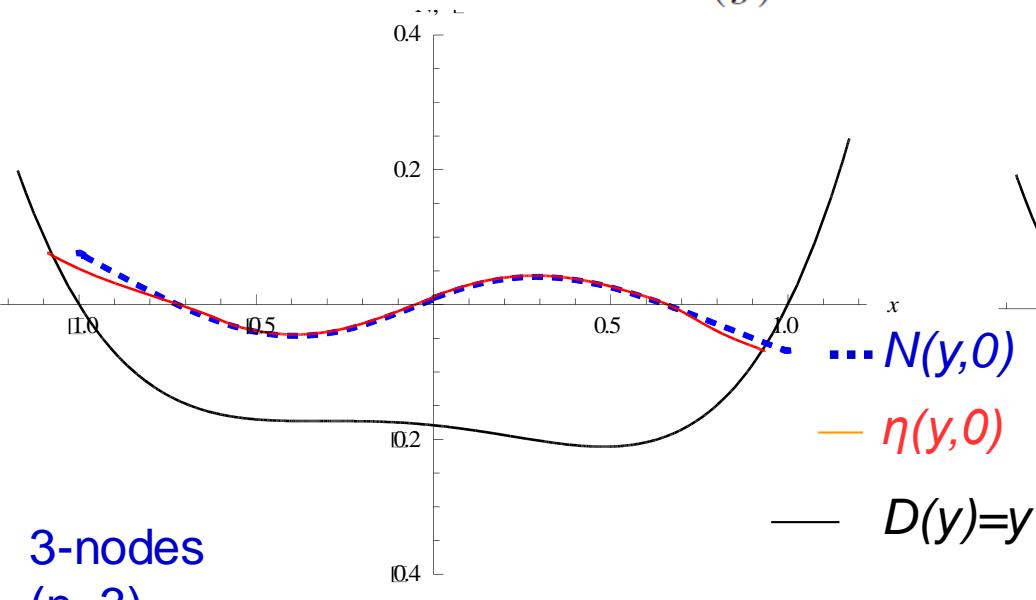
Reduced Carrier—Greenspan transform with cutting function ρ :

$$D(y) = D(x) + \eta(x, t)\rho(x), \quad \tau = t, \quad N(y, \tau) = \eta(x, t), \quad U(y, \tau) = u(x, t)$$

The leading term is defined from linearized shallow water with 2 fixed boundaries:

$$N_\tau + (D(y)U)_y = 0, \quad U_\tau + N_y = 0, \quad y \in [a, b], \quad E^2 = \|N\|^2 + \|\sqrt{y} U\|^2 < \infty$$

Finally: $x = y - \varepsilon N_1 \frac{\rho(y)}{D'(y)}$, $\eta(x, t) = N(y, t)$, $u(x, t) = U(y, t)$



Carrier-Greenspan transformation for standing waves and experimental studies

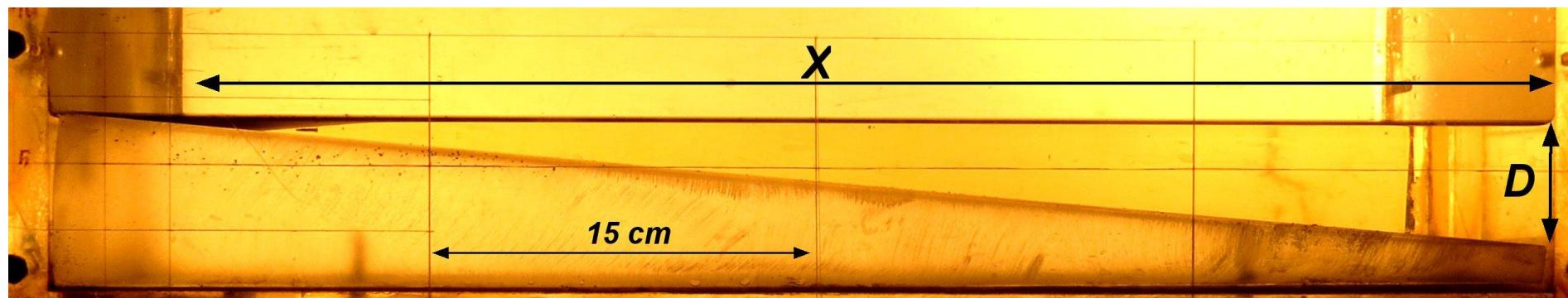
Nonlinear Shallow water equations:

$$\eta_t + ((D(x) + \eta)u)_x = 0, \quad u_t + g\eta_x + uu_x = 0, \quad D = \gamma(x - x^0), \quad u|_{x=b} = 0.$$

Carrier-Greenspan transformation:

The linear equation with nonlinear boundary condition

$$N_\tau + (yU)_y = 0, \quad U_\tau + N_y = 0, \quad U(Y(\tau), \tau) = 0, \quad Y(\tau) = b + N(Y(\tau), \tau), \\ \|N\|^2 + \|\sqrt{y} U\|^2 < \infty$$



Shallow water: sloping bottom, formal asymptotics

Formal series: $N(y, \tau, \varepsilon) = \varepsilon N_1^{w(\varepsilon)} + \varepsilon^2 N_2^{w(\varepsilon)} + \dots,$

Leading term: $N_1^{w(\varepsilon)} \equiv \cos(w_0\tau) \mathbf{J}_0(2w(\varepsilon)\sqrt{y}) \quad w_0 = \mu_n/2$
 $U_1^{w(\varepsilon)} \equiv \sin(w_0\tau) \frac{1}{\sqrt{y}} \mathbf{J}_1(2w(\varepsilon)\sqrt{y})$

Corrected frequency
to avoid resonances: $w(\varepsilon) = w_0 + \varepsilon^2 w_2 + \dots$

Boundary condition: $U_1(b, \tau) = 0$

$$U_2(b, \tau) = -U_{1y}(b, \tau)N_1(b, \tau) = \xi_1 \sin(2w_0\tau)$$

First correction: $N_2 = c_2 \cos(2w\tau) \mathbf{J}_0(4w\sqrt{y}) \quad c_2 = -\frac{w_0}{2} \mathbf{J}_0(2w_0) \mathbf{J}'_1(2w_0) (\mathbf{J}_1(4w_0))^{-1}$

Boundary for U_3 : $U_3^{w_0}(b, \tau) = \xi_2 \sin(3w_0\tau) + \xi_3 \sin(w_0\tau) - 2w_2 \mathbf{J}'_1(2w_0) \sin(w_0\tau)$

defines phase shift $w_2 = \xi_3/2\mathbf{J}'_1(2w_0)$

w_2 :

Second correction: $N_3 = c_3 \cos(3w_0\tau) \mathbf{J}_0(6w_0\sqrt{y}) \quad c_3 = \frac{\xi_2}{\mathbf{J}_1(6w_0)}$

Etc...

Shallow water: sloping bottom, the leading term

Formal series: $N(y, \tau, \varepsilon) = \varepsilon N_1^{w(\varepsilon)} + \varepsilon^2 N_2^{w(\varepsilon)} + \dots,$

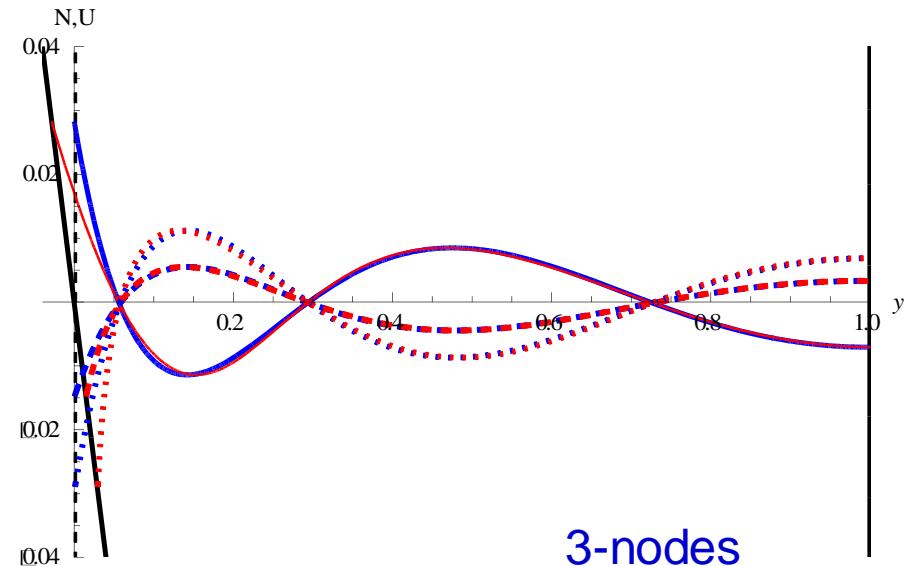
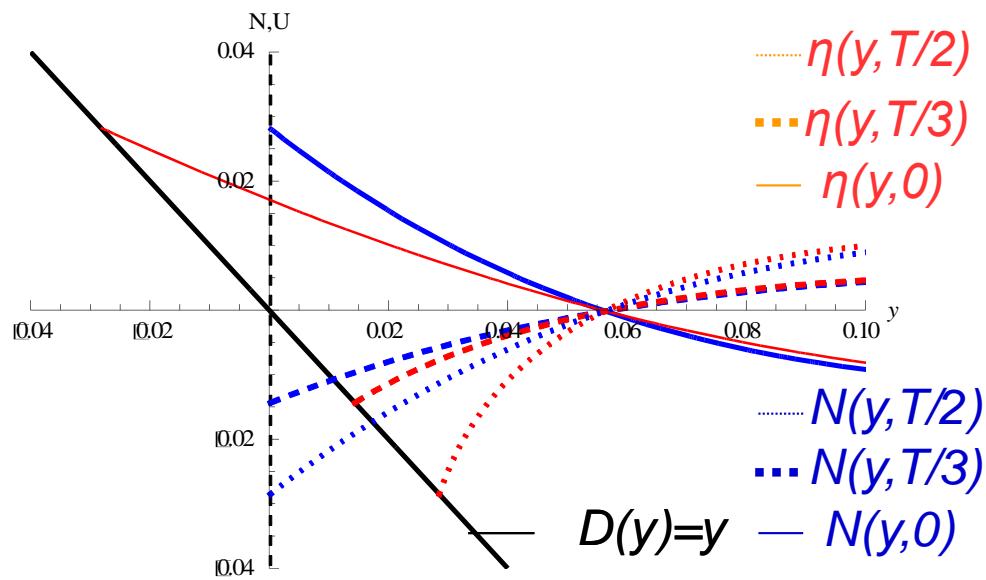
The Leading term
for linearized system: $N_1^{w(\varepsilon)} \equiv \cos(w_0 \tau) \mathbf{J}_0(2w(\varepsilon)\sqrt{y})$ $w_0 = \mu_n/2$
 $U_1^{w(\varepsilon)} \equiv \sin(w_0 \tau) \frac{1}{\sqrt{y}} \mathbf{J}_1(2w(\varepsilon)\sqrt{y})$

Reduced C—G transform: substitute $\tau(t, y, \varepsilon) = t + O(\varepsilon)$ into $N(y, \tau), U(y, \tau)$

And get the Leading term for shallow water –
parametrically defined via $y \in [0, 1]$:

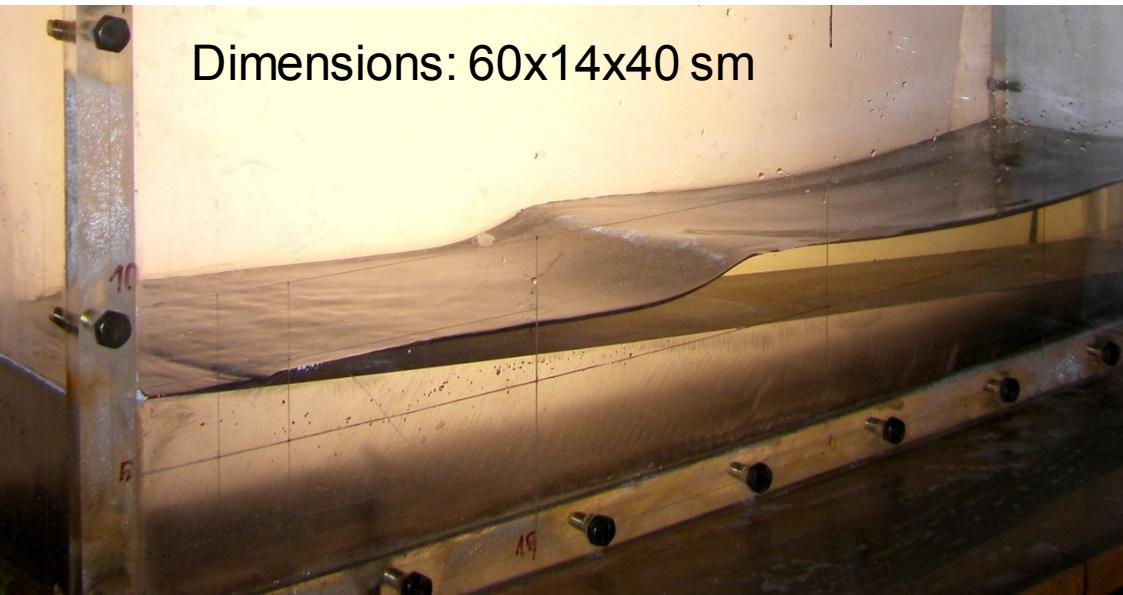
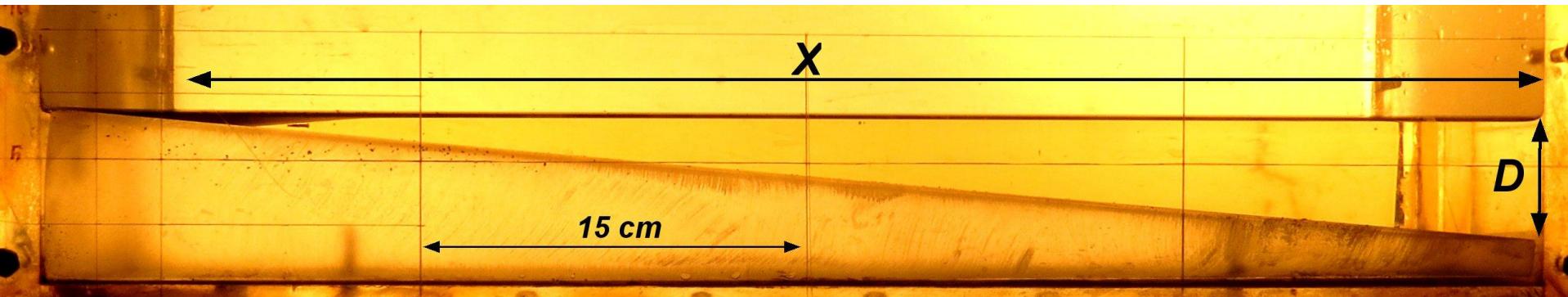
$$x = y - \varepsilon N_1^{w_0}(y, t) + O(\varepsilon^2)$$

$$u(x, t) = \varepsilon U_1^{w_0}(y, t) + O(\varepsilon^2)$$



3-nodes
(α, β)

Experimental setup: parametric resonance

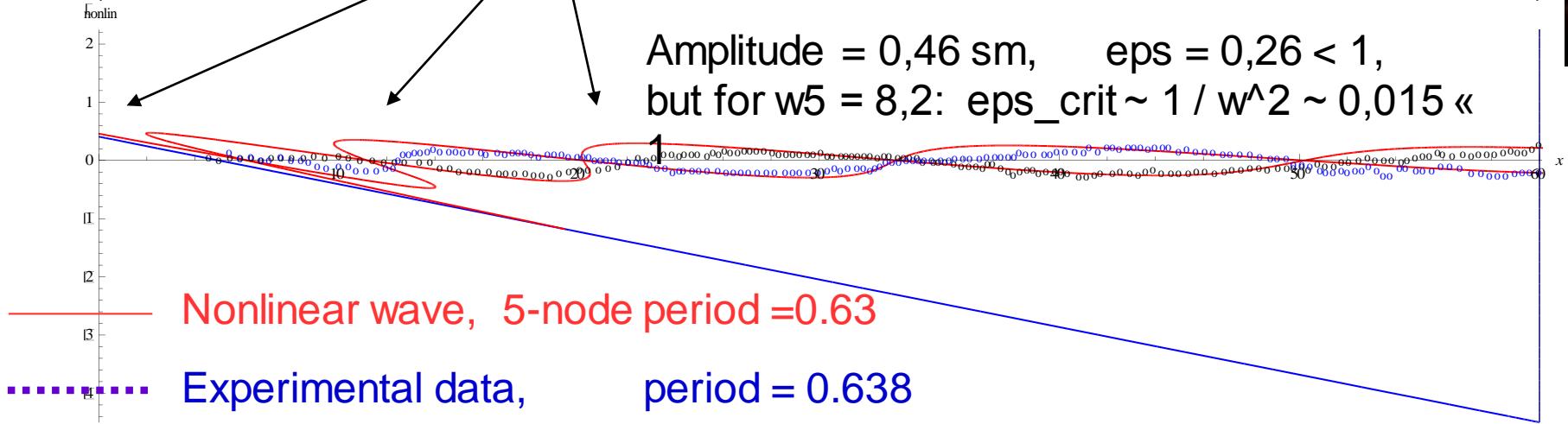
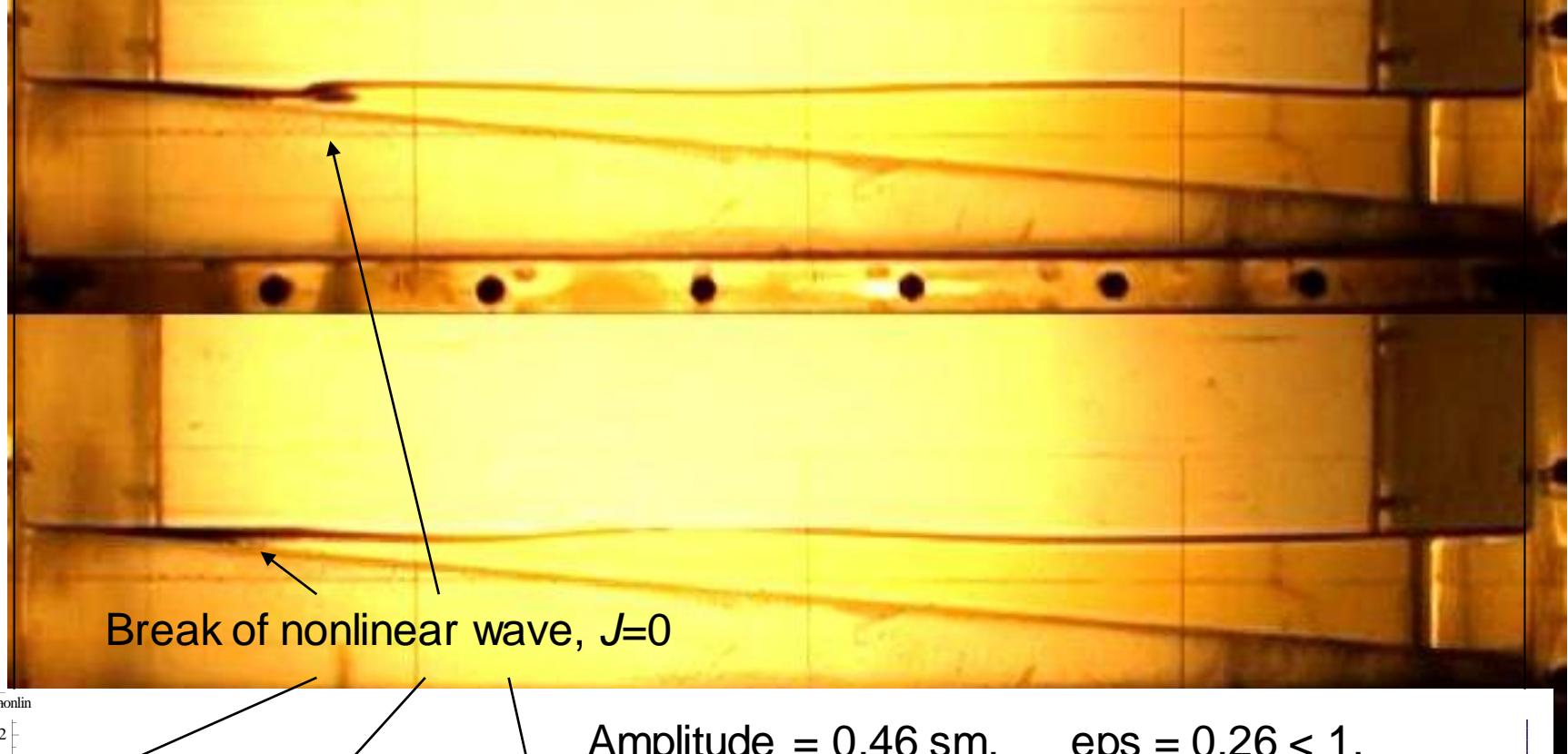


Gravity waves on the free surface
in rectangular vessel
(length = 60sm, width = 14 sm)
with slopping bottom ($D:X = 4,5\text{sm} : 55\text{sm}$)

Surface waves are induced
by vertical oscillations of vessel
with parametric resonance
(Oscillations period = waves period / 2)

Video 30 and 120 frames per sec,
editing in ImageJ

Experiment: comparison



Спасибо за внимание!

Будьте здоровы!