



# **Complex interactions of wave breathers**

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# Nonlinear Schrödinger Equation (NLSE) and its exact solutions

$$i\psi_t + \frac{1}{2}\psi_{xx} + |\psi|^2\psi = 0$$

The NLSE is integrable using the Inverse Scattering Transform (IST). The IST is based on the auxiliary Zakharov-Shabat linear system for the  $2 \times 2$  matrix wave function  $\Phi$ :

$$\Phi_x - \begin{pmatrix} -i\lambda & \psi \\ -\psi^* & i\lambda \end{pmatrix} \Phi = 0$$

$$\Phi_t - \begin{pmatrix} -i\lambda^2 + i|\psi|^2/2 & \lambda\psi + i\psi_x/2 \\ -\lambda\psi^* + i\psi_x^*/2 & i\lambda^2 - i|\psi|^2/2 \end{pmatrix} \Phi = 0$$

$\lambda$  is a complex-valued spectral parameter

Compatibility condition:

$$\Phi_{xt} = \Phi_{tx}$$

## Solitons in a parametrically unstable plasma

E. A. Kuznetsov

*Institute of Automation and Electrometry, Siberian Branch, Academy of Sciences of the USSR, Novosibirsk*

(Presented by Academician R. Z. Sagdeev, June 23, 1977)

(Submitted May 2, 1977)

Dokl. Akad. Nauk SSSR 236, 575-577 (September 1977)

PACS numbers: 52.35.Py, 52.35.Ra

$$i\psi_t + 1/2\psi_{xx} + (|\psi|^2 - |E_0|^2)\psi = 0$$

with the boundary condition

$$\psi \rightarrow E_0 \quad \text{as} \quad |x| \rightarrow \infty.$$

$$\xi = (\lambda^2 + |E_0|^2)^{1/2}$$

$$\lambda_{1,2} = \pm i\mu; \quad \xi_{1,2} = i\nu;$$

$$\Omega = 2\mu\nu.$$

$$E(x, t) = 2 \frac{\nu}{\mu} \left( \frac{\nu \cos \Omega t + i\mu \sin \Omega t}{\text{ch } 2\nu x + E_0/\mu \cos \Omega t} \right),$$

# Eigenvalue problem

$$\hat{L} = i \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix} \frac{\partial}{\partial x} - i \begin{pmatrix} 0 & \psi \\ \psi^* & 0 \end{pmatrix}$$

$$\hat{L}\Phi = \lambda\Phi,$$

Stable and unstable  
continuous spectrum

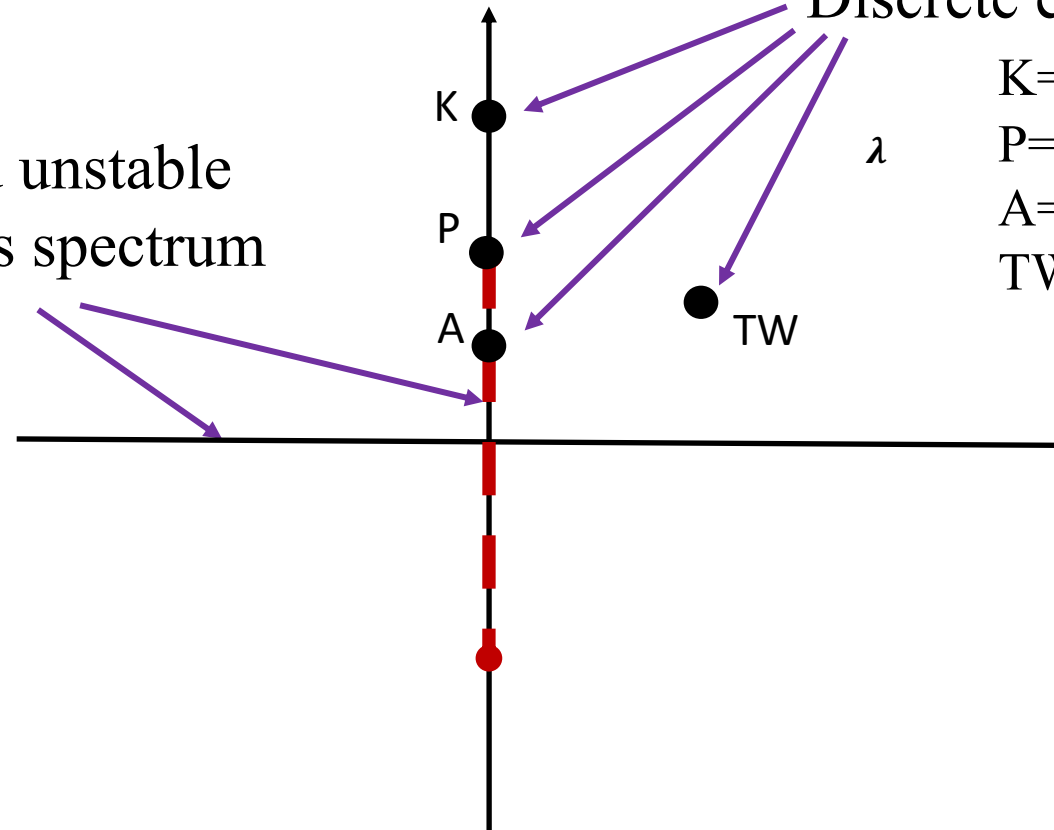
Discrete eigenvalues:

K=Kuznetsov

P=Peregrine

A=Akhmediev

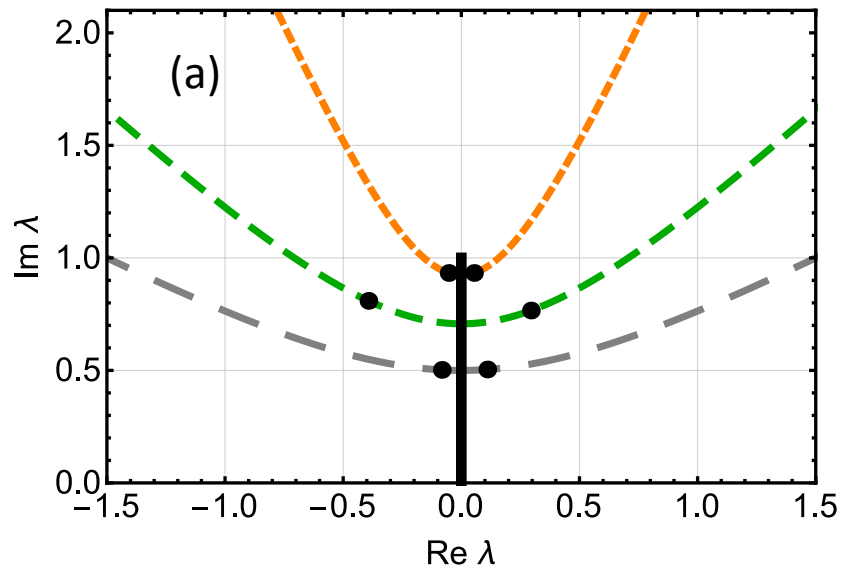
TW=Tajiri-Watanabe



$$\zeta = \sqrt{\lambda^2 + 1}$$

Branch cut  $[-i, i]$

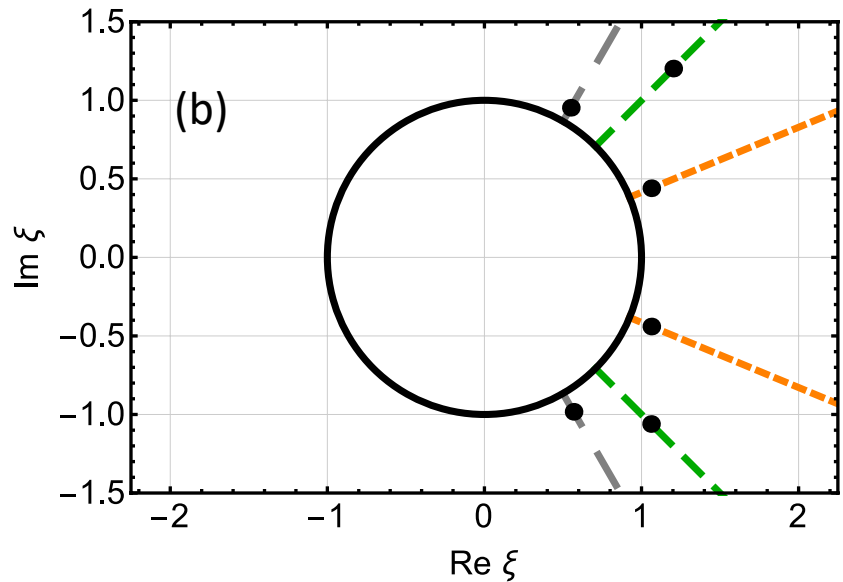
# Two different parametrizations for eigenvalues



$$\lambda = \frac{i}{2} \left( \xi + \frac{1}{\xi} \right), \quad \xi = R e^{i\alpha}$$

$$\operatorname{Re}[\lambda] = \pm \frac{\sin \alpha}{2} \left( R - \frac{1}{R} \right) \quad \operatorname{Im}[\lambda] = \frac{\cos \alpha}{2} \left( R + \frac{1}{R} \right) \quad (*)$$

$$\alpha_1 = -\alpha_2 \quad (**)$$



Comparison of  $\lambda$  and  $\xi$  parametrizations of the spectral parameter. The branch cut and its Joukowski mapping are drawn by black solid lines. The pairs of breather eigenvalues (marked by black points) lie on (a) the parametric curves Eq. (\*\*) (dashed lines) and (b) the rays Eq. (\*) (dashed lines).

## General $N$ -breather solution formula

$$\psi_N = e^{it} \left[ 1 + 2 \det \begin{pmatrix} 0 & q_{1,2} & \cdots & q_{N,2} \\ q_{1,1}^* & & & \\ \vdots & \widehat{M}^T & & \\ q_{N,1}^* & & & \end{pmatrix} (\det \widehat{M})^{-1} \right]$$

$$q_{i1} = e^{-\phi_i} - \frac{e^{\phi_i - i\alpha_i}}{R_i} \quad q_{i2} = e^{\phi_i} - \frac{e^{-\phi_i - i\alpha_i}}{R_i}$$

$$\phi_i = \eta_i(x - x_{0,i}) + \gamma_i t + i \left( k_i x + \delta_i t - \frac{\theta_i}{2} \right)$$

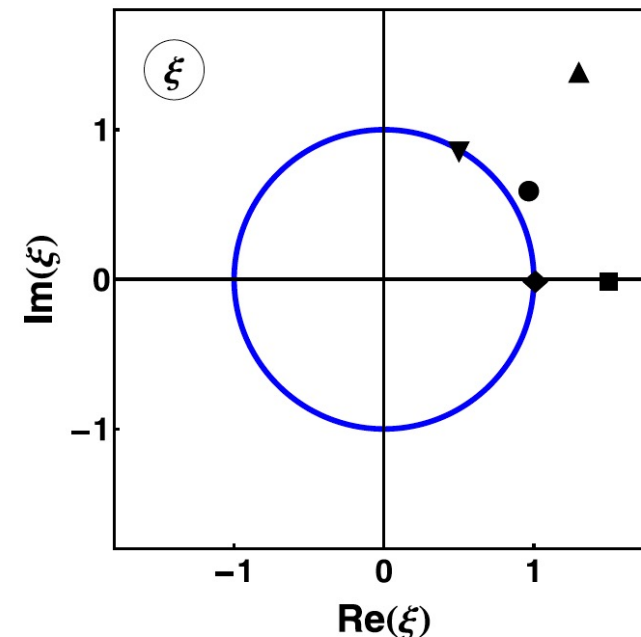
$$\eta_i = -\frac{1}{2} \left( R_i - \frac{1}{R_i} \right) \cos \alpha_i \quad \gamma_i = -\frac{1}{4} \left( R_i^2 + \frac{1}{R_i^2} \right) \sin 2\alpha_i$$

$$k_i = -\frac{1}{2} \left( R_i + \frac{1}{R_i} \right) \sin \alpha_i \quad \delta_i = \frac{1}{4} \left( R_i^2 - \frac{1}{R_i^2} \right) \cos 2\alpha_i$$

One breather has four real-valued parameters:  
 $R_i, \alpha_i$  - breather amplitude and group velocity  
 $x_{0,i}, \theta_i$  - position and phase

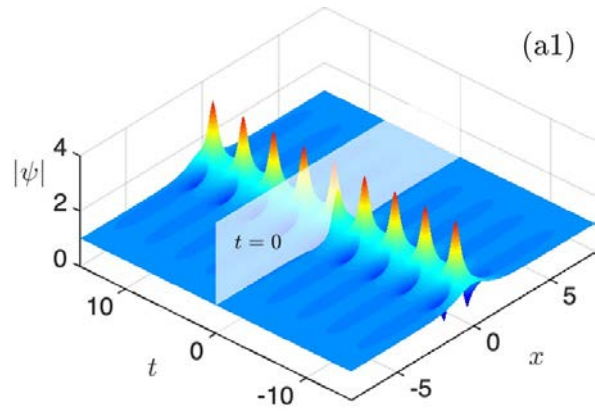
$$\widehat{M}_{nm} = \frac{i(\mathbf{q}_n \cdot \mathbf{q}_m^*)}{R_n e^{i\alpha_n} + \frac{e^{-i\alpha_n}}{R_n} - R_m e^{-i\alpha_m} - \frac{e^{i\alpha_m}}{R_m}}$$

$$\xi = R e^{i\alpha}$$

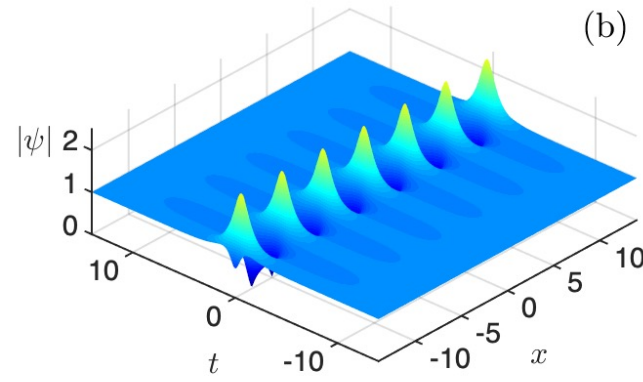


- ▲ - Tajiri-Watanabe,
- ▼ - Akhmediev breather, ● - quasi-Akhmediev breather
- - Kuznetsov breather, ◆ - Peregrine breather

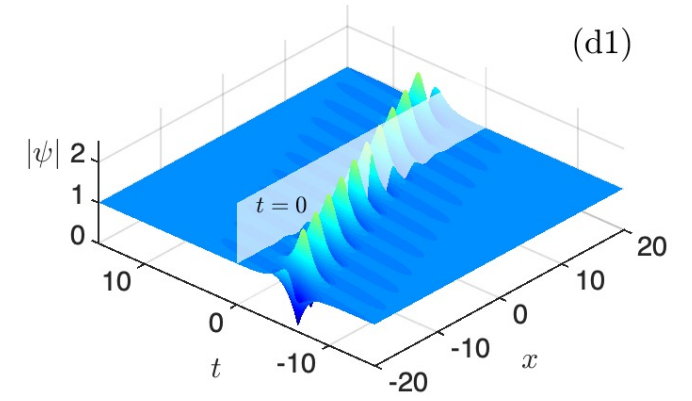
# Fundamental one-breather solutions



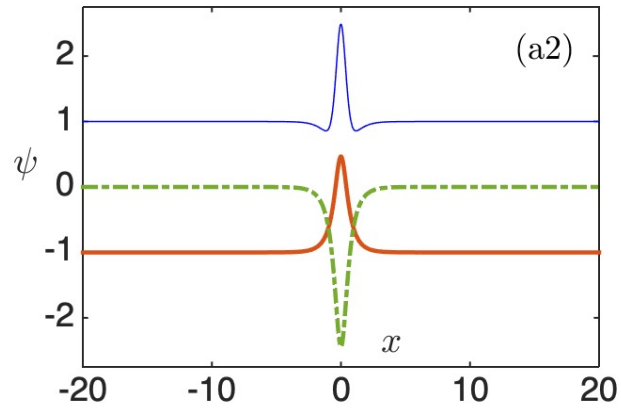
(a1)



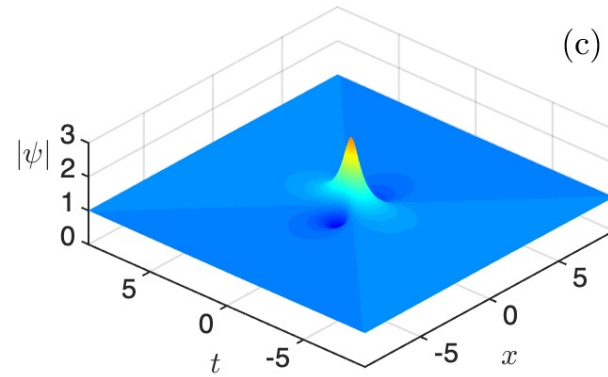
(b)



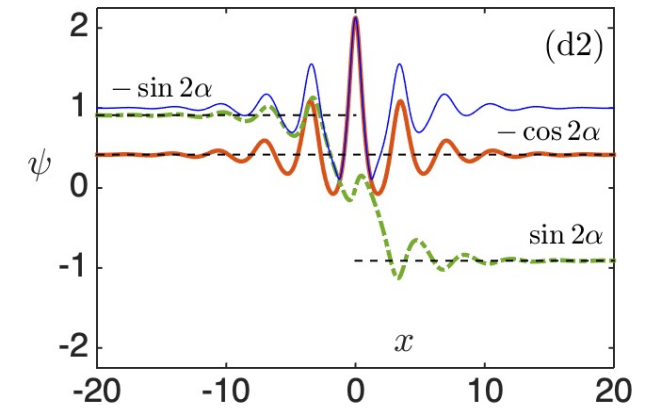
(d1)



(a2)



(c)



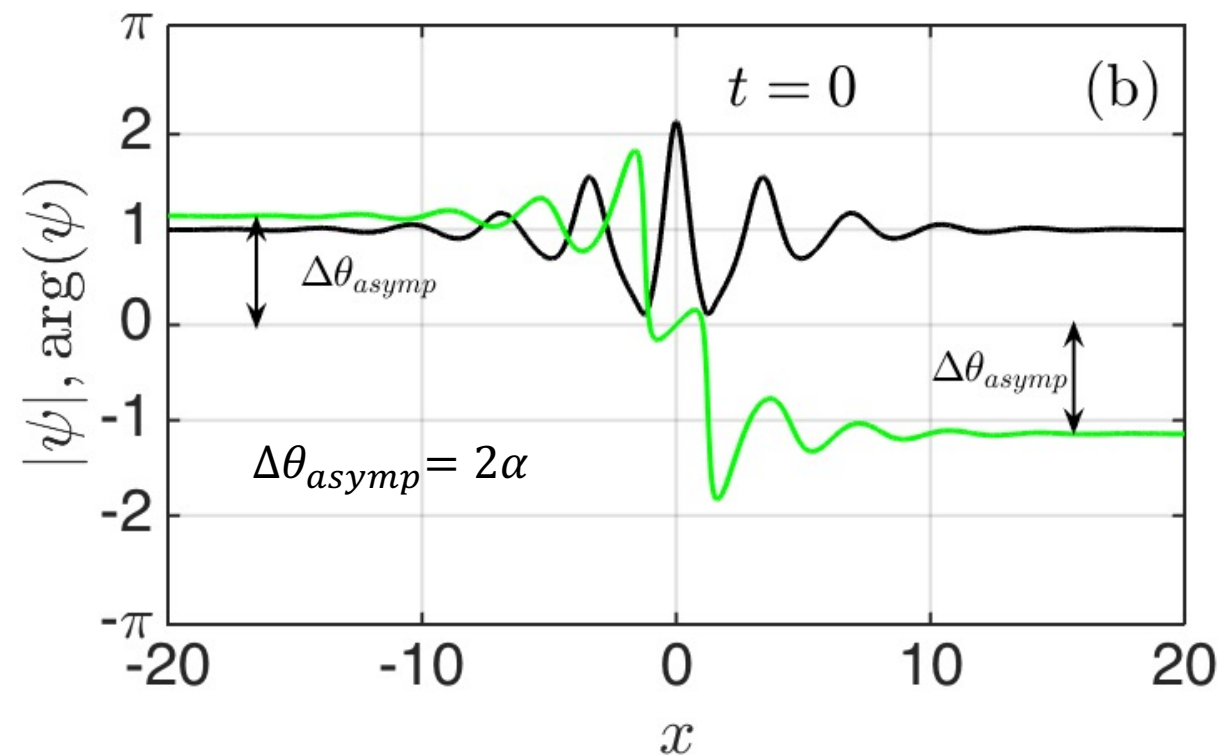
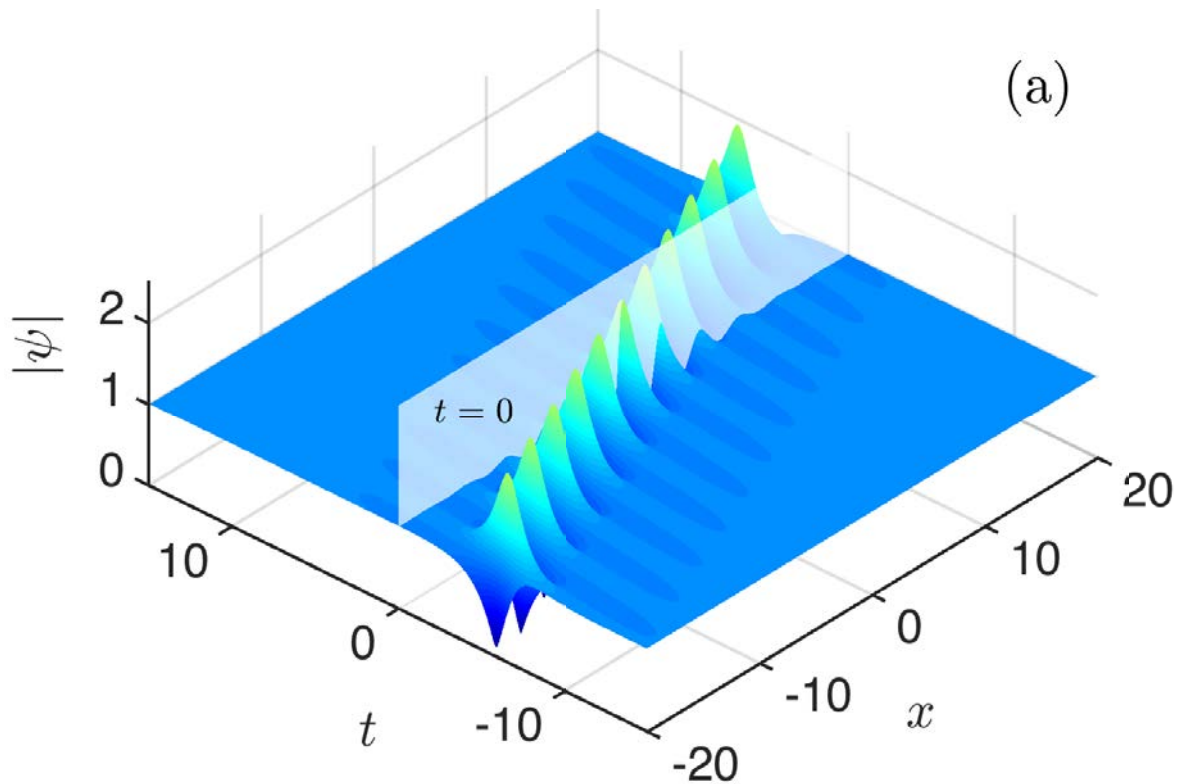
(d2)

(a) Kuznetsov breather

(b) Akhmediev breather  
(c) Peregrine breather

(d) Tajiri-Watanabe breather

## General single-breather solution: $R > 1, \alpha > 0$

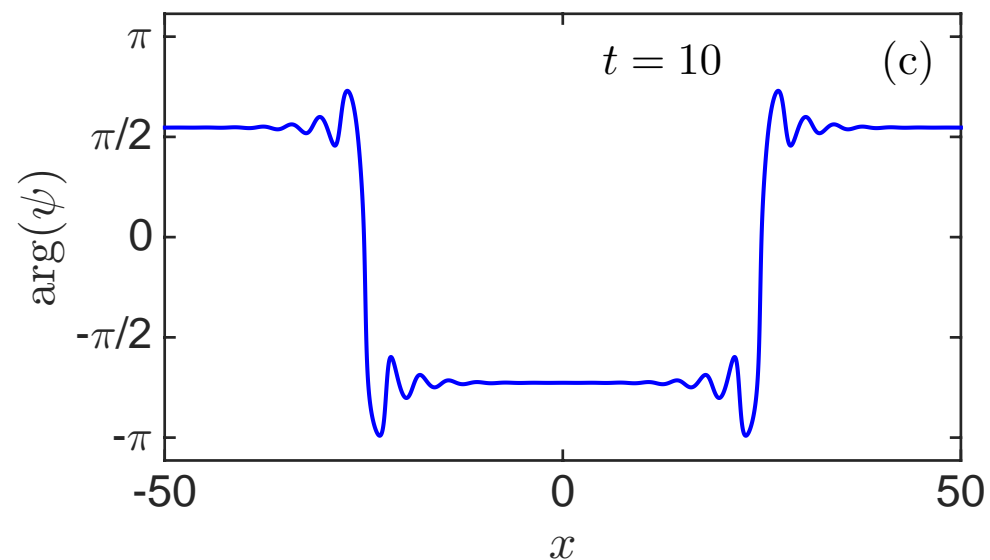
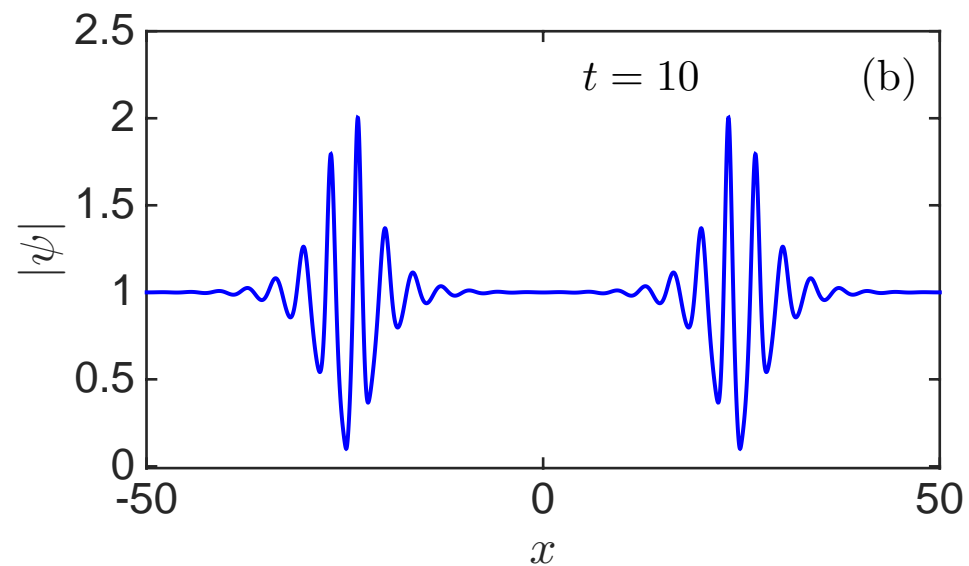
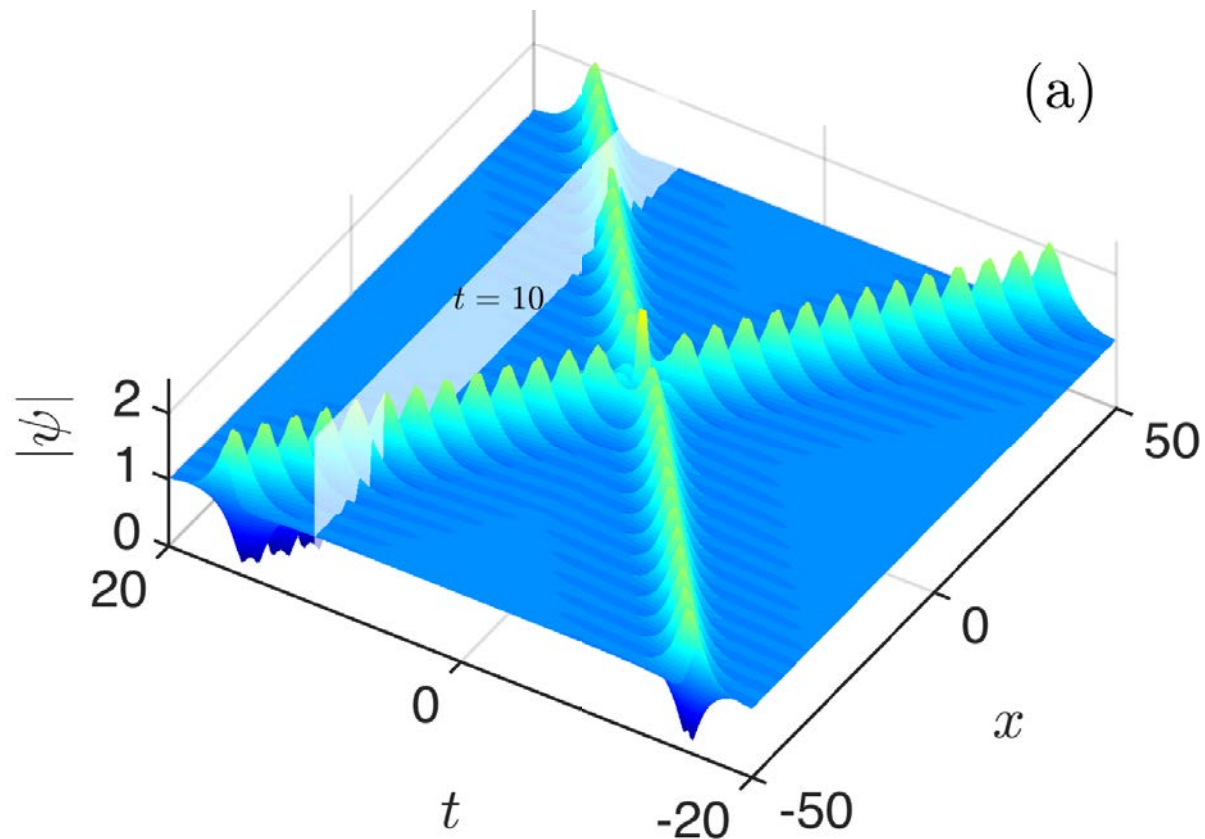


$$\psi = \left( 1 + 2\left(R + \frac{1}{R}\right) \cos \alpha \frac{q_1^* q_2}{|q_1|^2 + |q_2|^2} \right) e^{it} e^{i\theta_c}$$

$$V_{gr} = -\sin \alpha (R^4 + 1) / (R(R^2 - 1))$$

$$V_{ph} = -\frac{1}{2} \left( R^2 - \frac{1}{R^2} \right) \cos 2\alpha$$

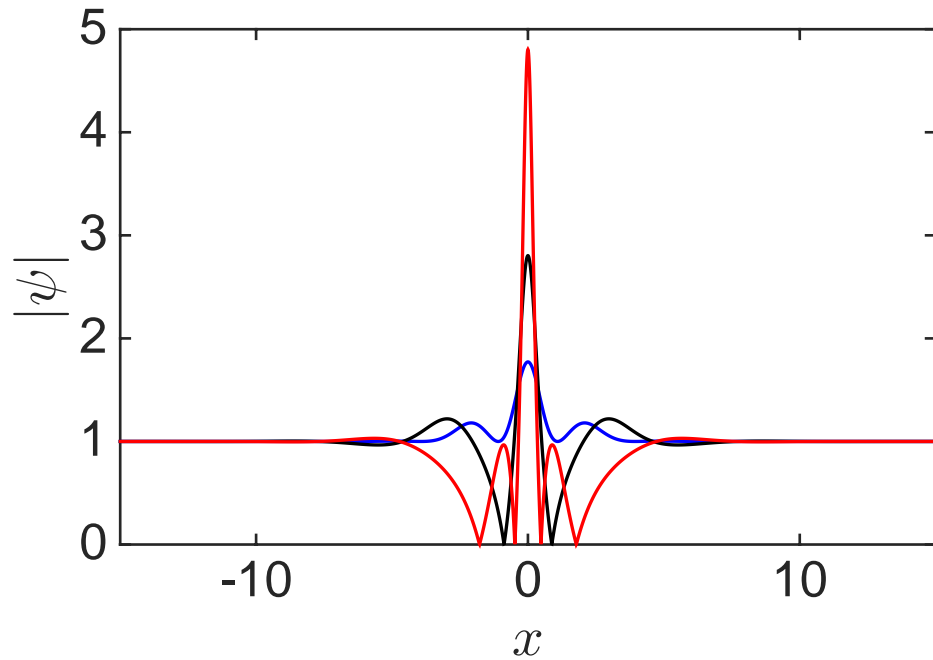
# Two-breather solution with $R_1 = R_2 = R$ ; $\alpha_1 = \alpha$ ; $\alpha_2 = -\alpha$ ; $R > 1, \alpha > 0$ .



Typical two-breather collision and its asymptotic state. (a) shows spatial-temporal portrait of a typical two-breather solution with  $R = 1.35$ ,  $\alpha = 1.0$ ,  $\theta_1 = 0$ ,  $\theta_2 = 0.025$ . Blue solid lines in (b) and (c) shows spatial profile of the solution at  $\xi = 10$ .

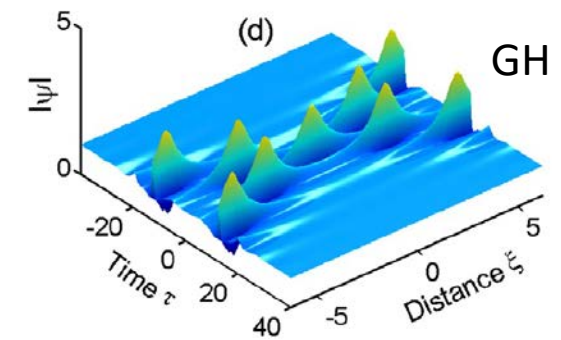
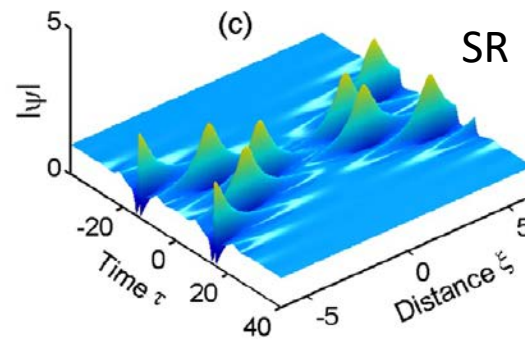
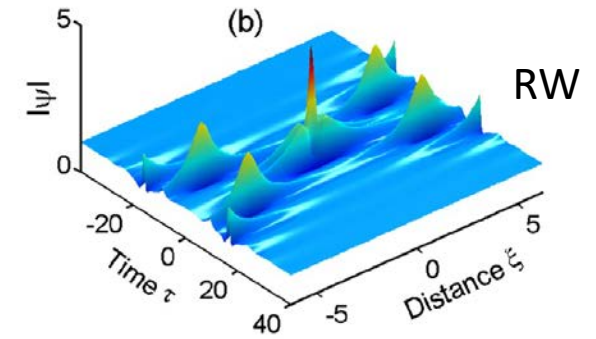
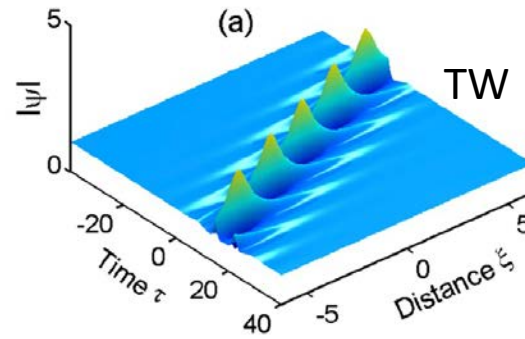


# Phase synchronization of two-breather collision



Amplitude profiles at  $t=0$  for the basic cases of two-breather collision: superregular (SR, blue line), rogue wave (RW, red line) and ghost (GH, black line) phase synchronizations.

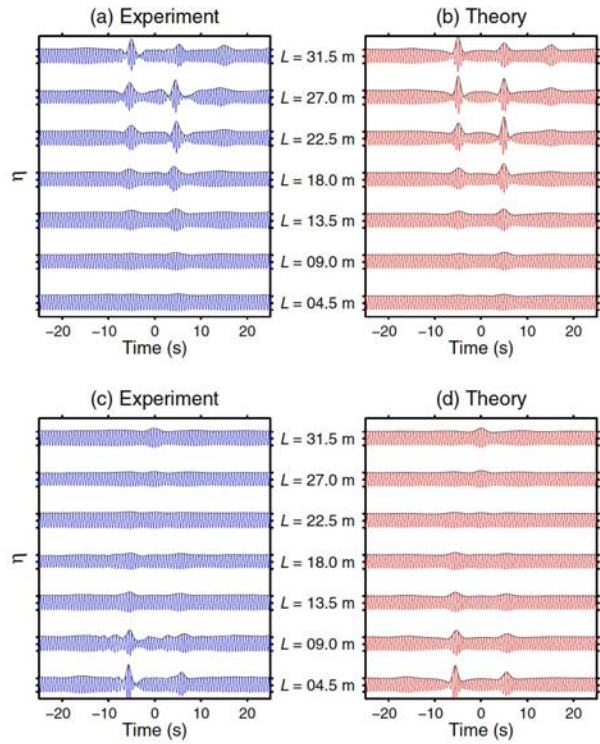
$$\theta_{\text{SR}} = \{\pi/2; 3\pi/2\}, \quad \theta_{\text{RW}} = 0, \quad \theta_{\text{GH}} = \pi.$$



[1,2]

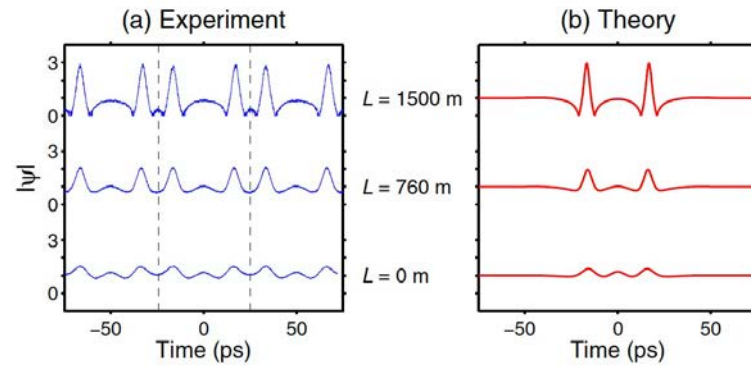
# Experimental observation of SR and GH breather interactions

## Water waves



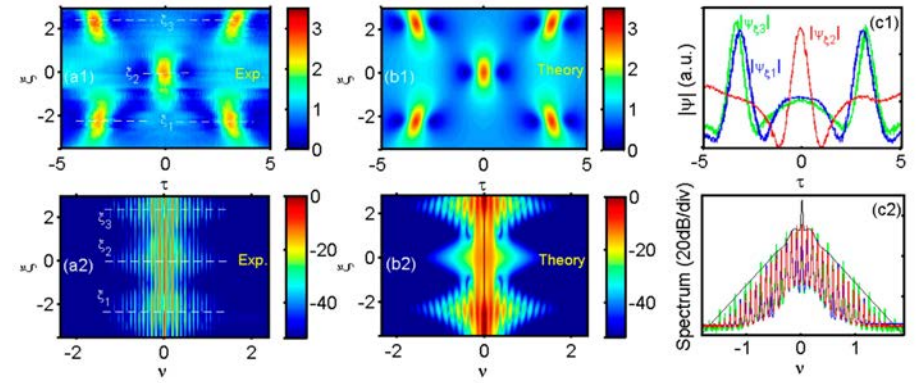
[2]

## Light waves

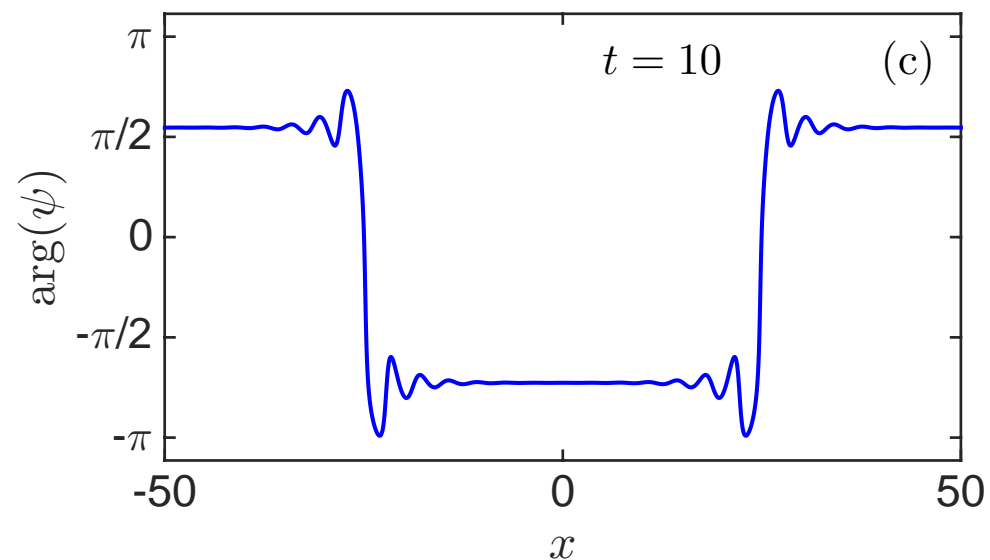
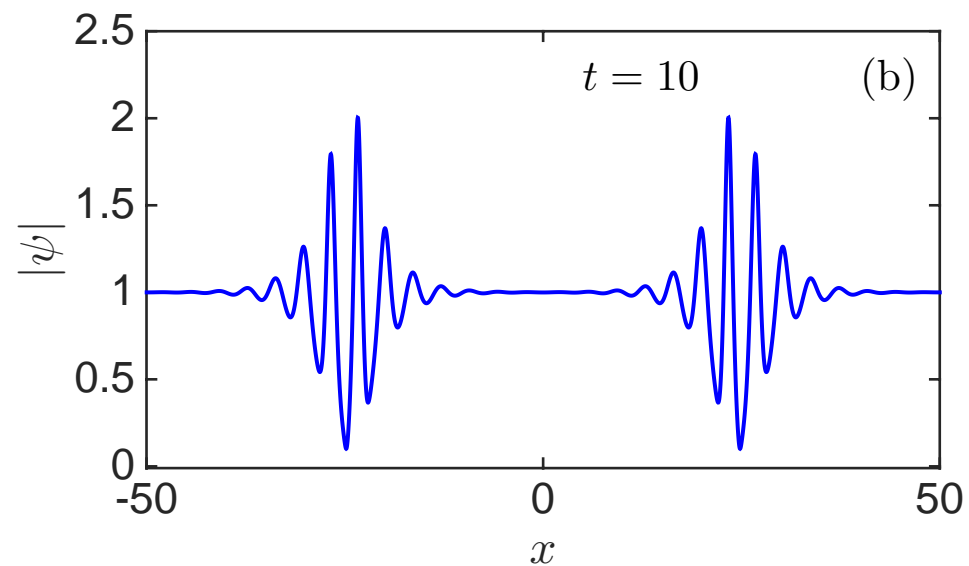
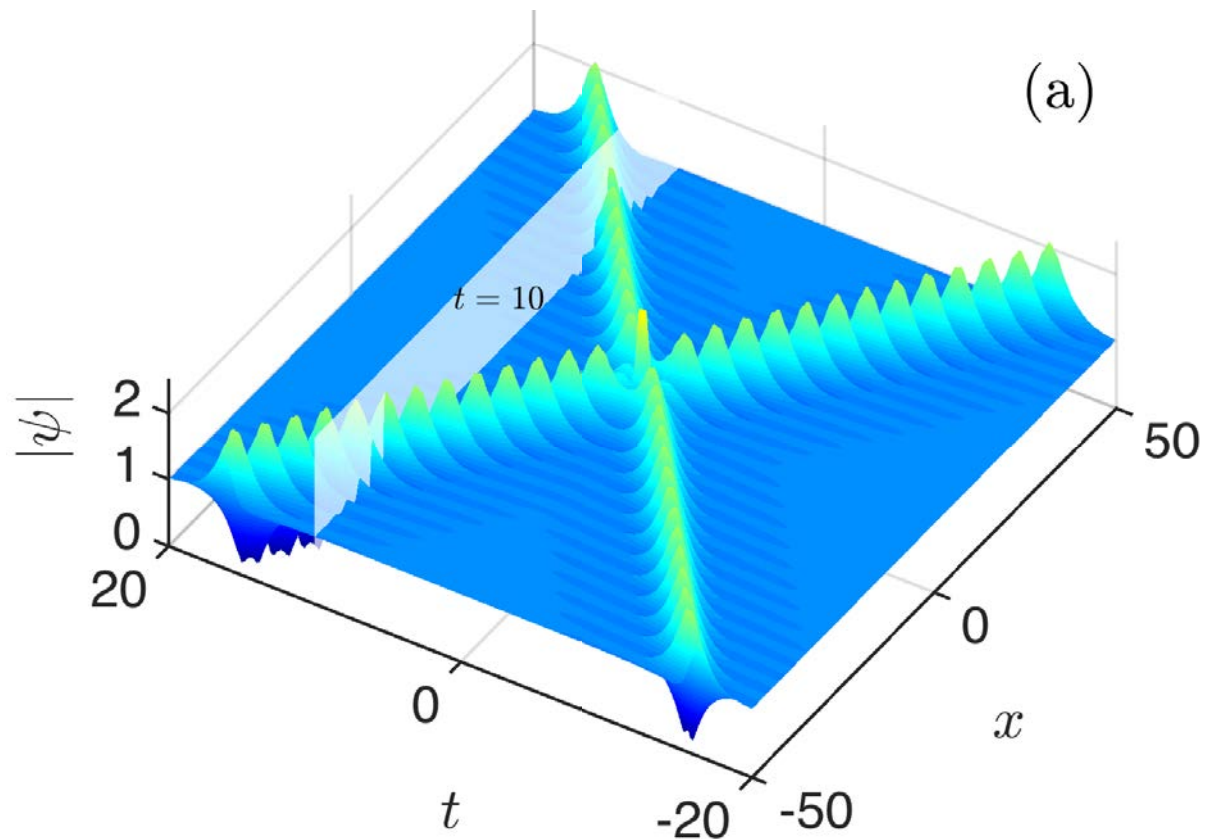


[3]

## Light waves



## Two-breather solution with $R_1 = R_2 = R$ ; $\alpha_1 = \alpha$ ; $\alpha_2 = -\alpha$ ; $R > 1, \alpha > 0$ .



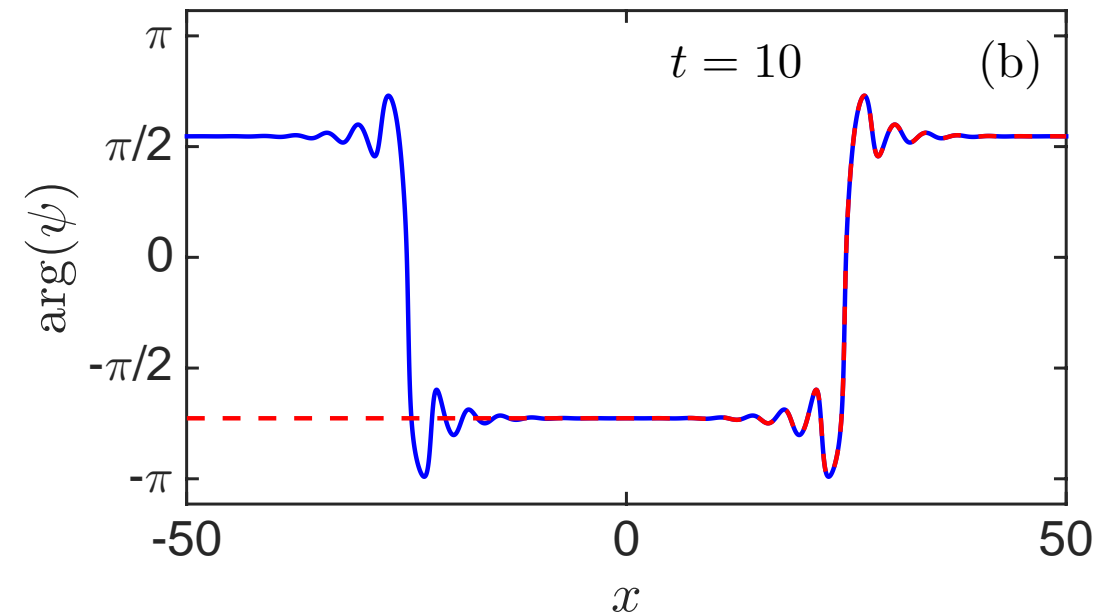
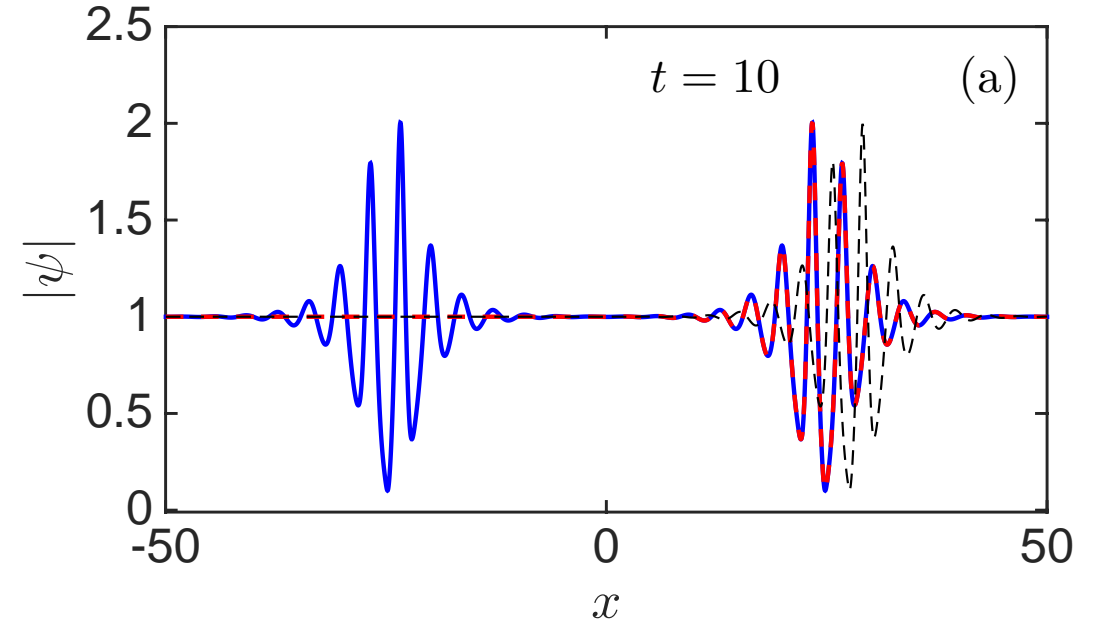
Typical two-breather collision and its asymptotic state. (a) shows spatial-temporal portrait of a typical two-breather solution with  $R = 1.35$ ,  $\alpha = 1.0$ ,  $\theta_1 = 0$ ,  $\theta_2 = 0.025$ . Blue solid lines in (b) and (c) shows spatial profile of the solution at  $\xi = 10$ .

# Space-phase shifts formulas [4]

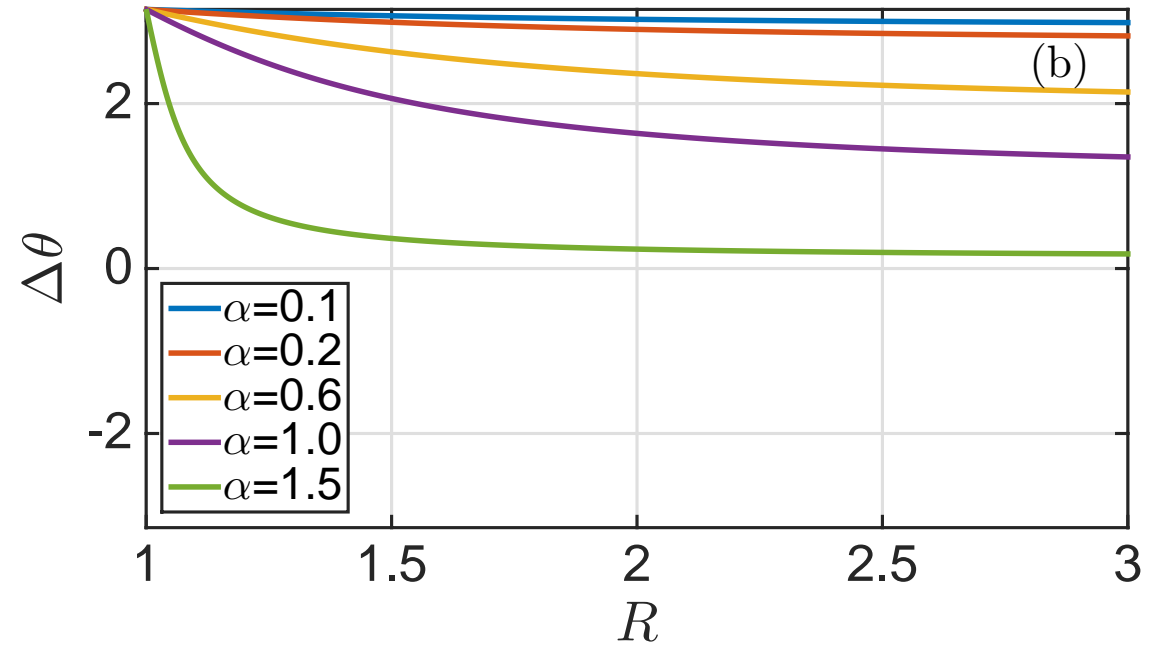
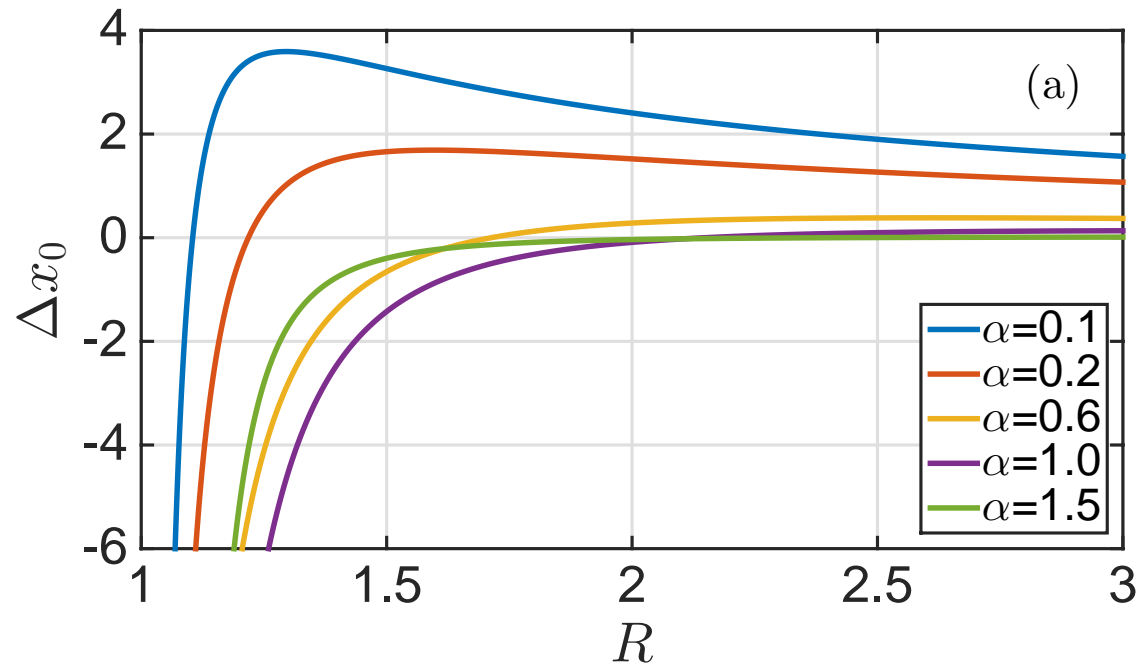
$$\Delta x_0 = \frac{\ln \left[ \frac{\left(R - \frac{1}{R}\right)^2}{\sin^2 \alpha \left(R^2 + \frac{1}{R^2} + 2 \cos 2\alpha\right)} \right]}{\left(R - \frac{1}{R}\right) \cos \alpha}$$

$$\Delta \theta_0 = -2 \arg \left[ \frac{2 R^3 (1 - R^2) \tan \alpha}{(1 + R^2)^2 \cos \alpha + i(1 - R^4)^2 \sin \alpha} \right]$$

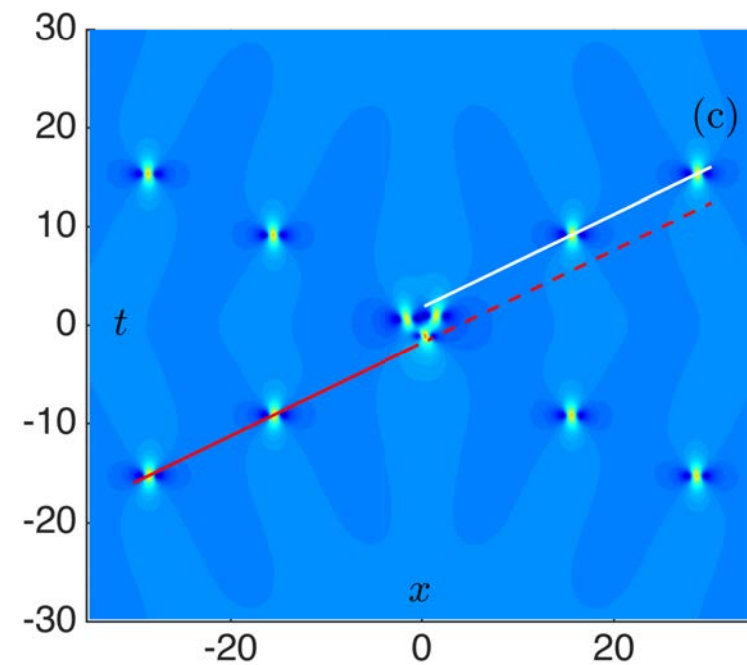
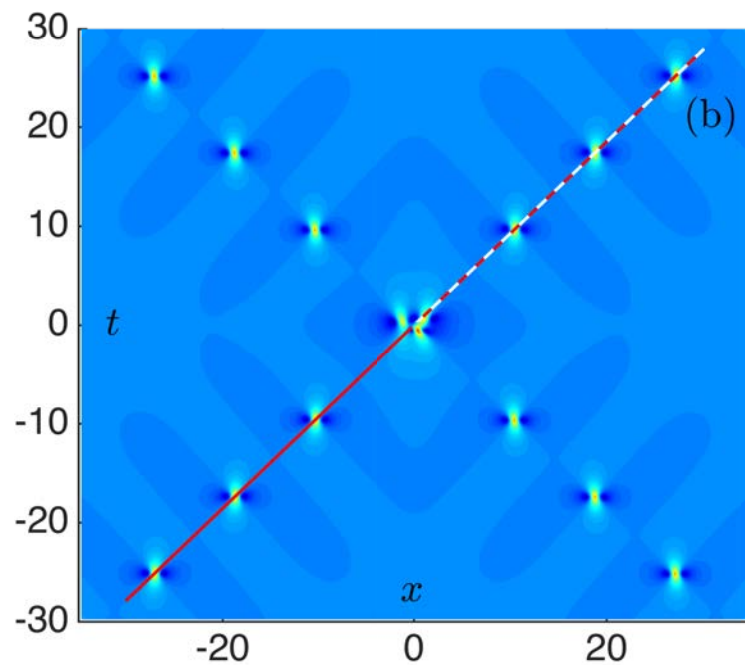
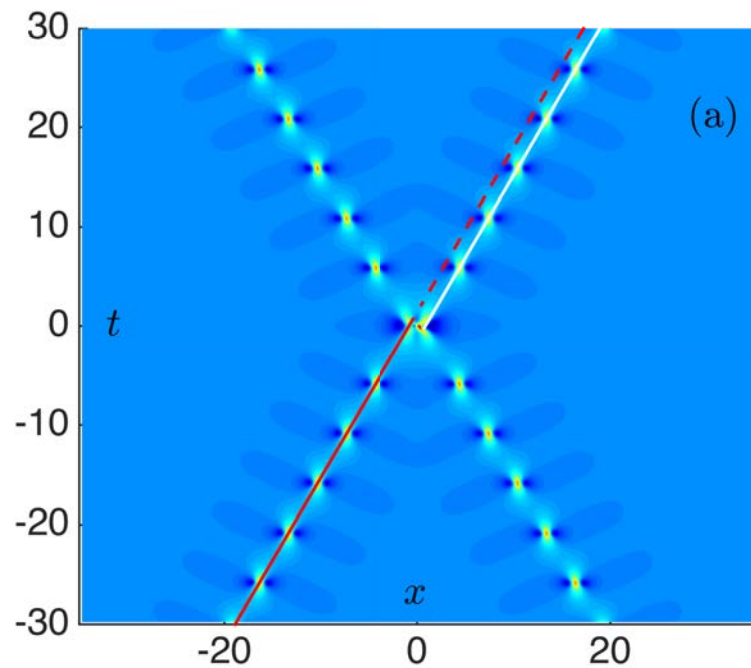
$$\Delta \theta_c = 4 \alpha$$



# Dependence of space and phase shifts on breather parameters



# Positive, zero and negative space shifts acquiring by breathers after collision.





## Space-phase shifts formulas in general case [5]

$$\Delta x_{0,21} = \ln [(s_1 - s_3)/(s_2 - s_4)] / 2 \operatorname{Im}[\zeta_2],$$

$$\Delta \theta_{21} = 2 \operatorname{Arg}[i(p_1 + p_2)], \Delta \theta_{c,21} = -4 \operatorname{Arg}[\lambda_1 + \zeta_1].$$

$$s_1 = (A^4 + |\lambda_1 + \zeta_1|^2 \cdot |\lambda_2 + \zeta_2|^2) \cdot |\lambda_1 - \lambda_2^*|^2 +$$

$$+ A^2 (|\lambda_1 + \zeta_1|^2 + |\lambda_2 + \zeta_2|^2) \cdot |\lambda_1 - \lambda_2|^2,$$

$$s_2 = A^2 (|\lambda_1 + \zeta_1|^2 + |\lambda_2 + \zeta_2|^2) \cdot |\lambda_1 - \lambda_2^*|^2 +$$

$$+ (A^4 + |\lambda_1 + \zeta_1|^2 \cdot |\lambda_2 + \zeta_2|^2) \cdot |\lambda_1 - \lambda_2|^2,$$

$$s_3 = A^2 (\lambda_1 - \lambda_1^*) (\lambda_2 - \lambda_2^*) \cdot [(\lambda_1 + \zeta_1)(\lambda_2 + \zeta_2) +$$

$$+ (\lambda_1^* + \zeta_1^*)(\lambda_2^* + \zeta_2^*)],$$

$$s_4 = -A^2 (\lambda_1 - \lambda_1^*) (\lambda_2 - \lambda_2^*) \cdot [(\lambda_1 + \zeta_1)(\lambda_2^* + \zeta_2^*) +$$

$$+ (\lambda_1^* + \zeta_1^*)(\lambda_2 + \zeta_2)];$$

$$p_1 = \{A^2 (\lambda_2 + \zeta_2 - \lambda_1^* - \zeta_1^*) - |\lambda_1 + \zeta_1|^2 (\lambda_2^* + \zeta_2^*) +$$

$$+ |\lambda_2 + \zeta_2|^2 (\lambda_1 + \zeta_1)\} / \{|\lambda_1 - \lambda_2^*|^2\},$$

$$p_2 = \frac{(A^2 + |\lambda_1 + \zeta_1|^2) (\lambda_2 + \zeta_2 - \lambda_2^* - \zeta_2^*)}{(\lambda_1 - \lambda_1^*) (\lambda_2 - \lambda_2^*)}.$$

**Well-known limit of solitons on zero background:**

$$\lim_{A \rightarrow 0} \Delta x_{0,21} = \frac{1}{2 \operatorname{Im}[\lambda_2]} \ln \left( \frac{|\lambda_1 - \lambda_2^*|^2}{|\lambda_1 - \lambda_2|^2} \right)$$

$$\lim_{A \rightarrow 0} \Delta \theta_{0,21} + \Delta \theta_{c,21} = -2 \operatorname{Arg} \left[ \frac{\lambda_1 - \lambda_2}{\lambda_1^* - \lambda_2} \right]$$

## Some details of the calculations

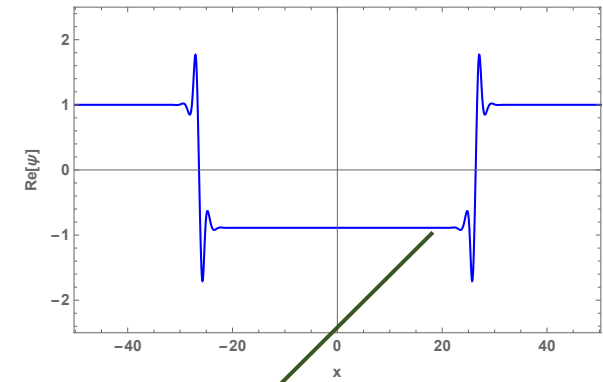
$$\psi_2 = A + 2 \frac{\det \begin{pmatrix} 0 & q_{1,2} & q_{2,2} \\ q_{1,1}^* & i \frac{|q_1|^2}{\lambda_1 - \lambda_1^*} & i \frac{(q_2 \cdot q_1^*)}{\lambda_2 - \lambda_1^*} \\ q_{1,N}^* & i \frac{(q_1 \cdot q_2^*)}{\lambda_1 - \lambda_2^*} & i \frac{|q_2|^2}{\lambda_2 - \lambda_2^*} \end{pmatrix}}{\det \begin{pmatrix} i \frac{|q_1|^2}{\lambda_1 - \lambda_1^*} & i \frac{(q_1 \cdot q_2^*)}{\lambda_1 - \lambda_2^*} \\ i \frac{(q_2 \cdot q_1^*)}{\lambda_2 - \lambda_1^*} & i \frac{|q_2|^2}{\lambda_2 - \lambda_2^*} \end{pmatrix}}$$

General two-breather solution

$$q_{n,1} = e^{-\phi_n} - \frac{iAe^{\phi_n}}{\lambda_n + \zeta_n}, \quad q_{n,2} = -\frac{iAe^{-\phi_n}}{\lambda_n + \zeta_n} + e^{\phi_n},$$

$$\phi_n = -i\zeta_n x - \text{Im}[\zeta_n]x_{0,n} - i\lambda_n \zeta_n t - i\theta_n/2$$

$$\zeta_n = \sqrt{\lambda_n^2 + A^2}$$



$$\psi_2 = e^{-i\theta_{2c}} \left( A + 2i(\lambda_2 + \lambda_2^*) \frac{\tilde{q}_{2,1}^* \tilde{q}_{2,2}}{|\tilde{q}_2|^2} \right)$$

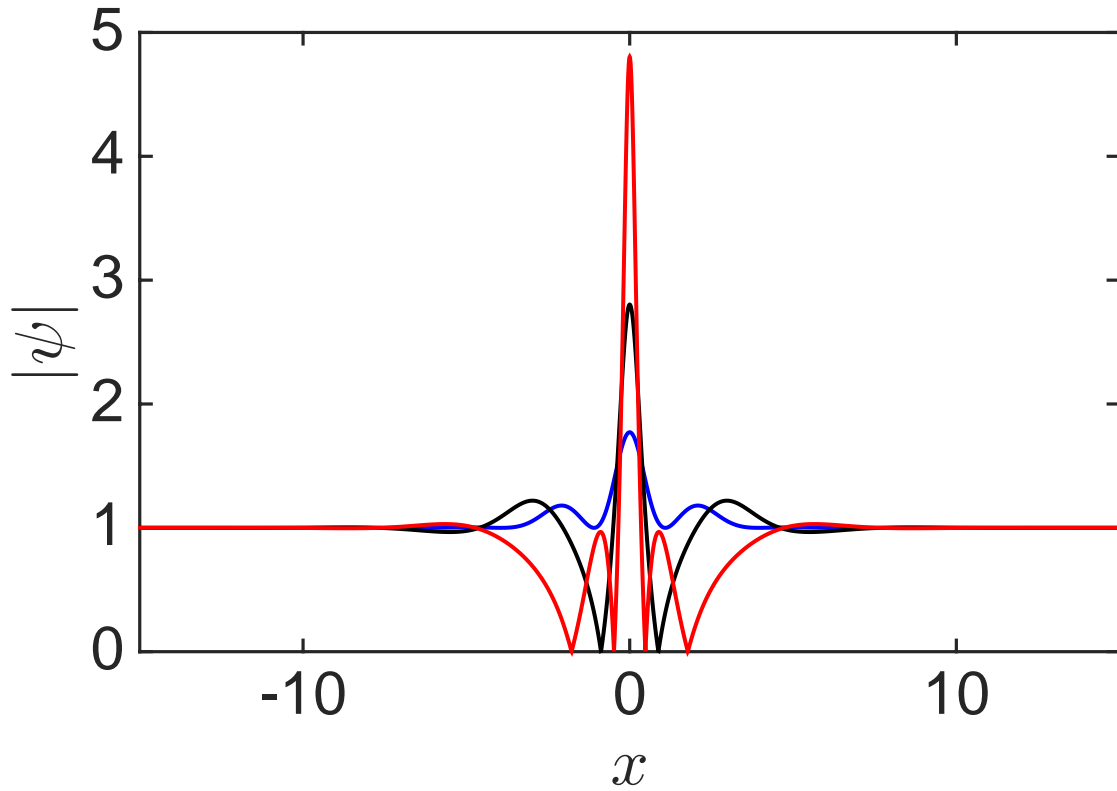
Asymptotic one-breather solution

$$t \rightarrow \infty, \quad x \sim t V_{gr2} \quad e^{\phi_1} \rightarrow \infty, \quad e^{-\phi_1} \rightarrow 0$$

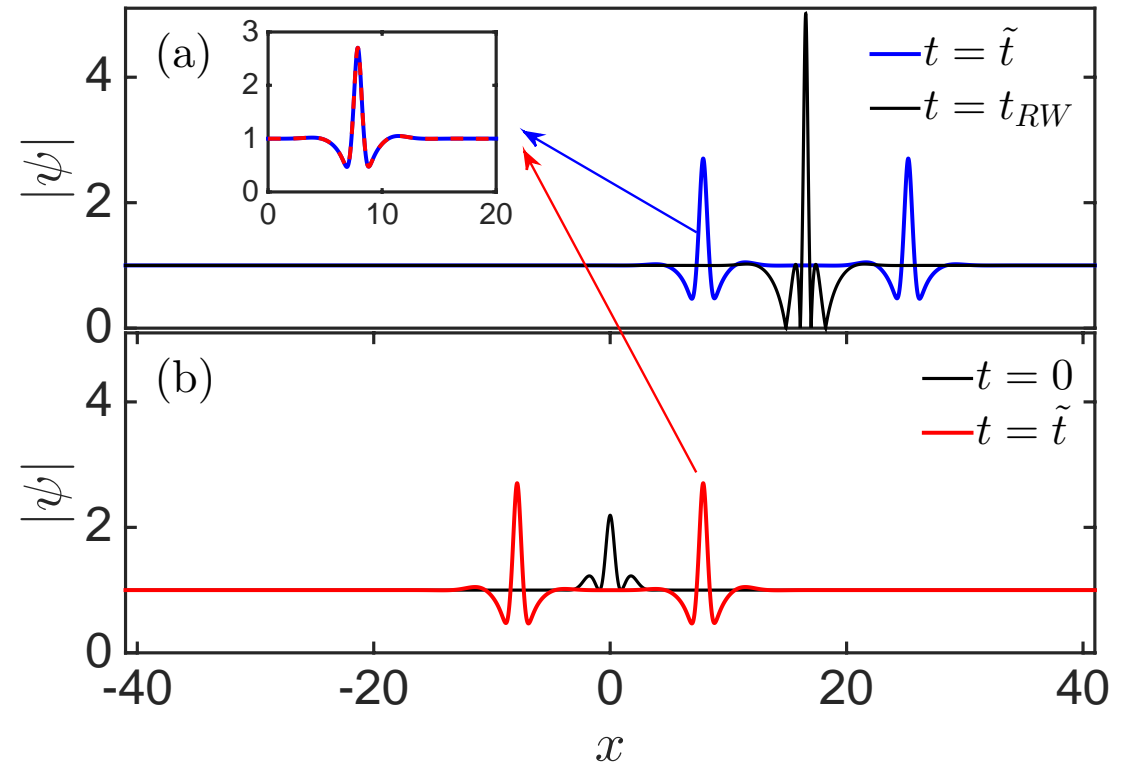
$$\begin{aligned} \tilde{q}_n(x, t, x_{0,n}, \theta_n) &= \\ &= q_n(x, t, x_{0,n} + \Delta x_{0,n}, \theta_n + \Delta \theta_n) \end{aligned}$$



# Main idea of synchronization of breather interactions



Amplitude profiles of symmetric two-breather collisions at  $t = 0$ . Blue – SR synchronization, RW – rogue wave synchronization.



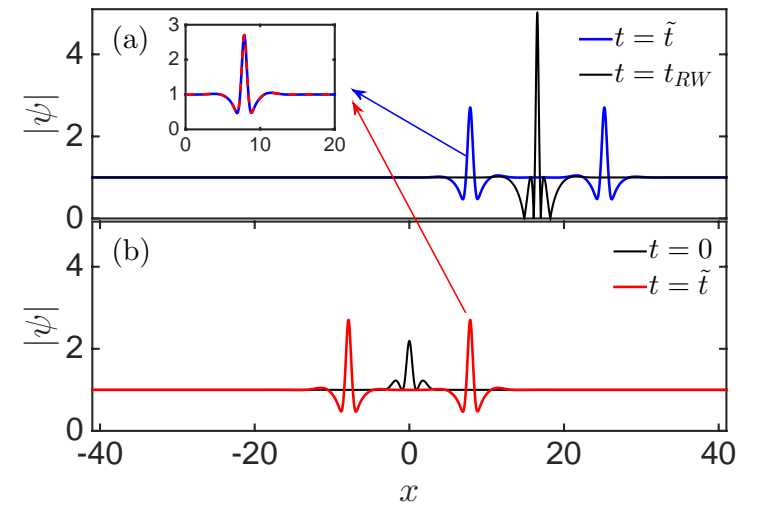
$$\phi = \eta(x - x_0) + \gamma t + i \left( kx + \delta t - \frac{\theta}{2} \right)$$

$$V_{gr} = -\frac{\gamma}{\eta} \quad V_{ph} = -2\delta$$

Rogue wave at  $\delta x_{synch}$  and  $t_{RW} = t_0$ :

$$\theta = -V_{ph} t_0 + 2k\delta x_{synch}$$

$$x_0 = -V_{gr} t_0 + \delta x_{synch}$$



$$\phi_{2Br}^{RW} = \eta(x + V_{gr} t_0 - \delta x_{synch}) + \gamma t + i \left( kx + \delta t + \frac{V_{gr} t_0 - 2k\delta x_{synch} + \theta_i}{2} \right)$$

$$\phi_{1Br}^{RW} = \phi_{2Br}^{RW}(x_0 \rightarrow x_0 - \Delta x_0, \theta \rightarrow \theta - \Delta\theta) = \eta(x + \Delta x_0 + V_{gr} t_0 - \delta x_{synch}) + \gamma t + i \left( kx + \delta t + \frac{V_{gr} t_0 + \Delta\theta - 2k\delta x_{synch}}{2} \right)$$

$$\phi_{2Br}^{SR} = \eta x + \gamma t + i \left( kx + \delta t - \frac{\pi/2}{2} \right)$$

$$\phi_{1Br}^{SR} = \phi_{2Br}^{SR}(x_0 \rightarrow x_0 + \Delta x_0, \theta \rightarrow \theta + \Delta\theta) = \eta(x - \Delta x_0) + \gamma t + i \left( kx + \delta t - \frac{\theta_f}{2} \right)$$

$$\phi_{1Br}^{RW} = \phi_{1Br}^{SR}$$

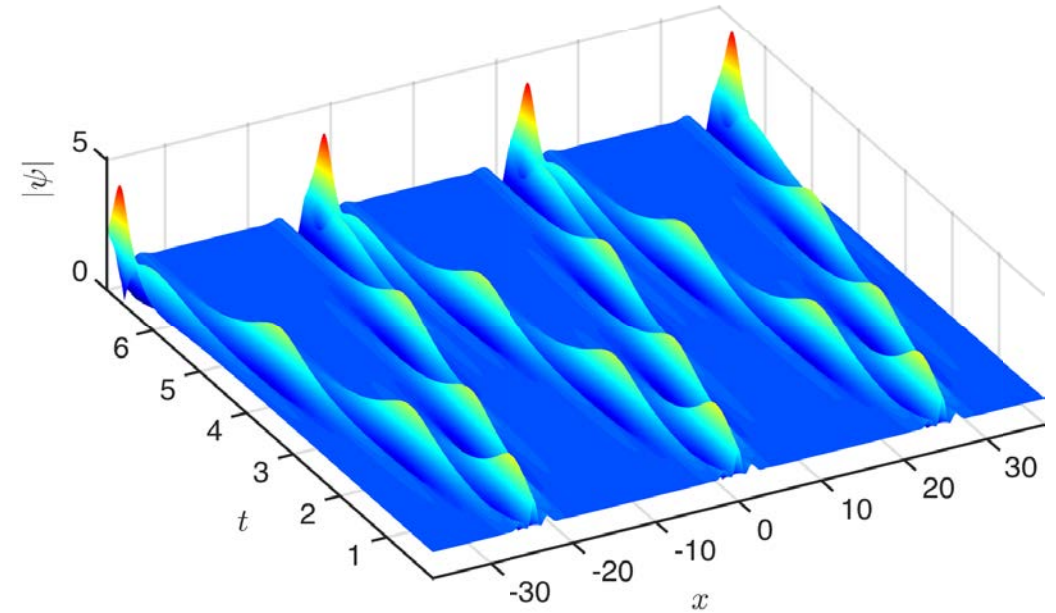
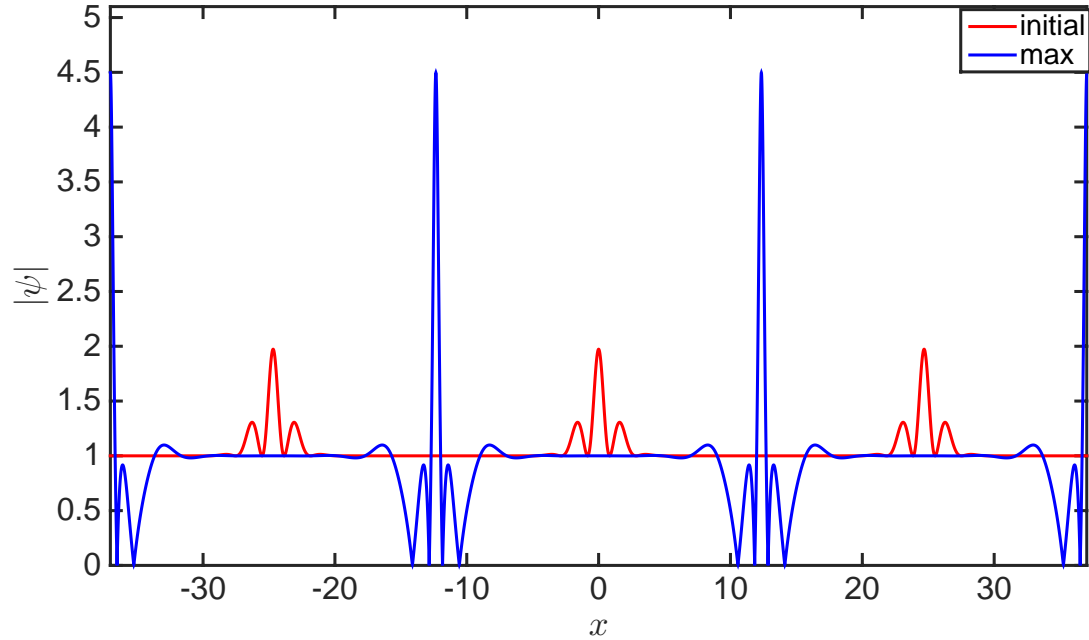


**Synchronization**



$$\delta x_{synch} = \frac{2\Delta x_0 + (2\Delta\theta_0 + \theta_i + \theta_f + 2\pi n) \frac{V_{gr}}{V_{ph}}}{1 - 2k \frac{V_{gr}}{V_{ph}}}$$

# Synchronization: periodic train of SR perturbations



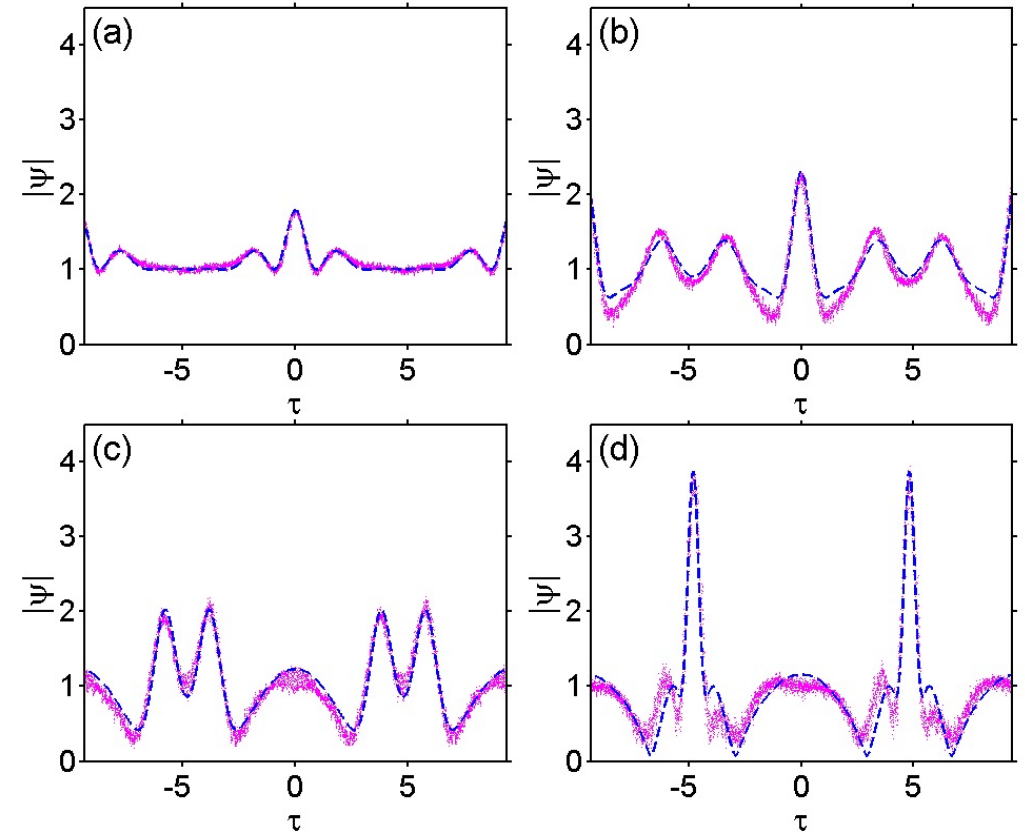
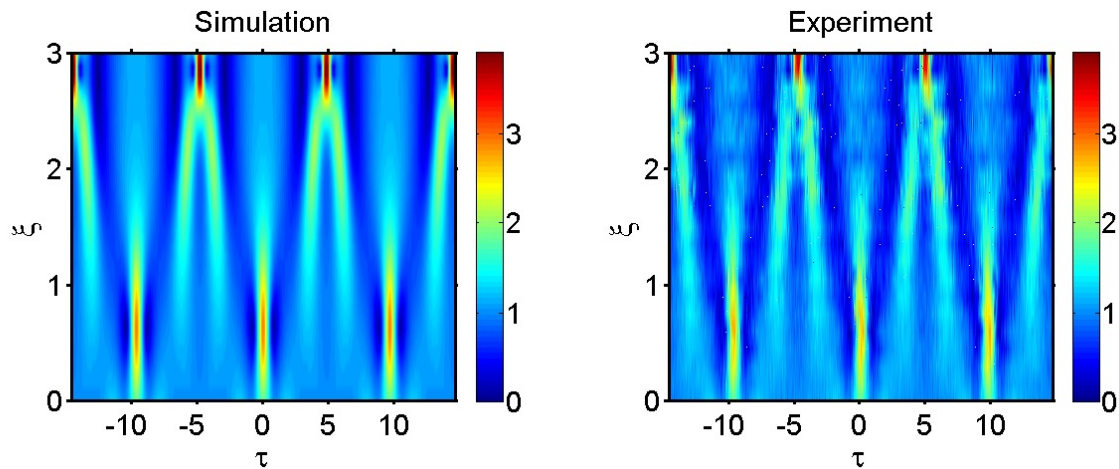
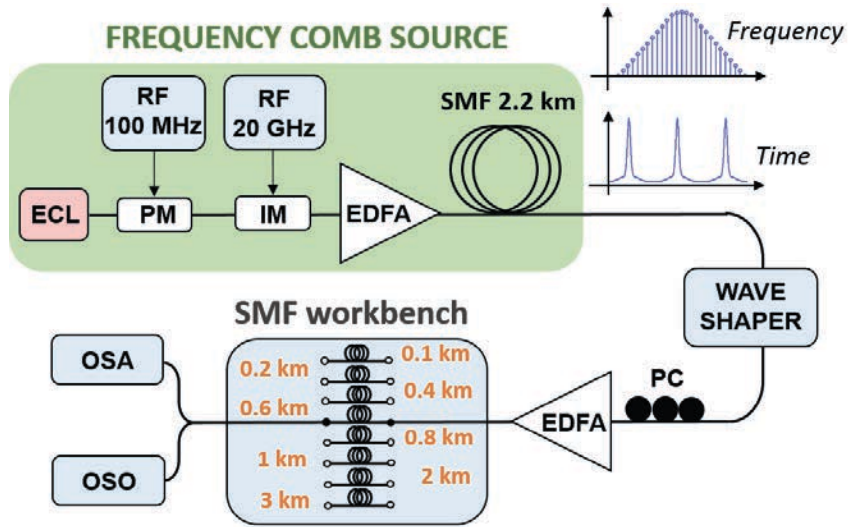
$$V_{gr} = V_{gr_2} = \sin \alpha (R^4 + 1) / (R(R^2 - 1))$$

$$V_{ph} = V_{ph_2} = -\frac{1}{2} \left( R^2 - \frac{1}{R^2} \right) \cos 2\alpha$$

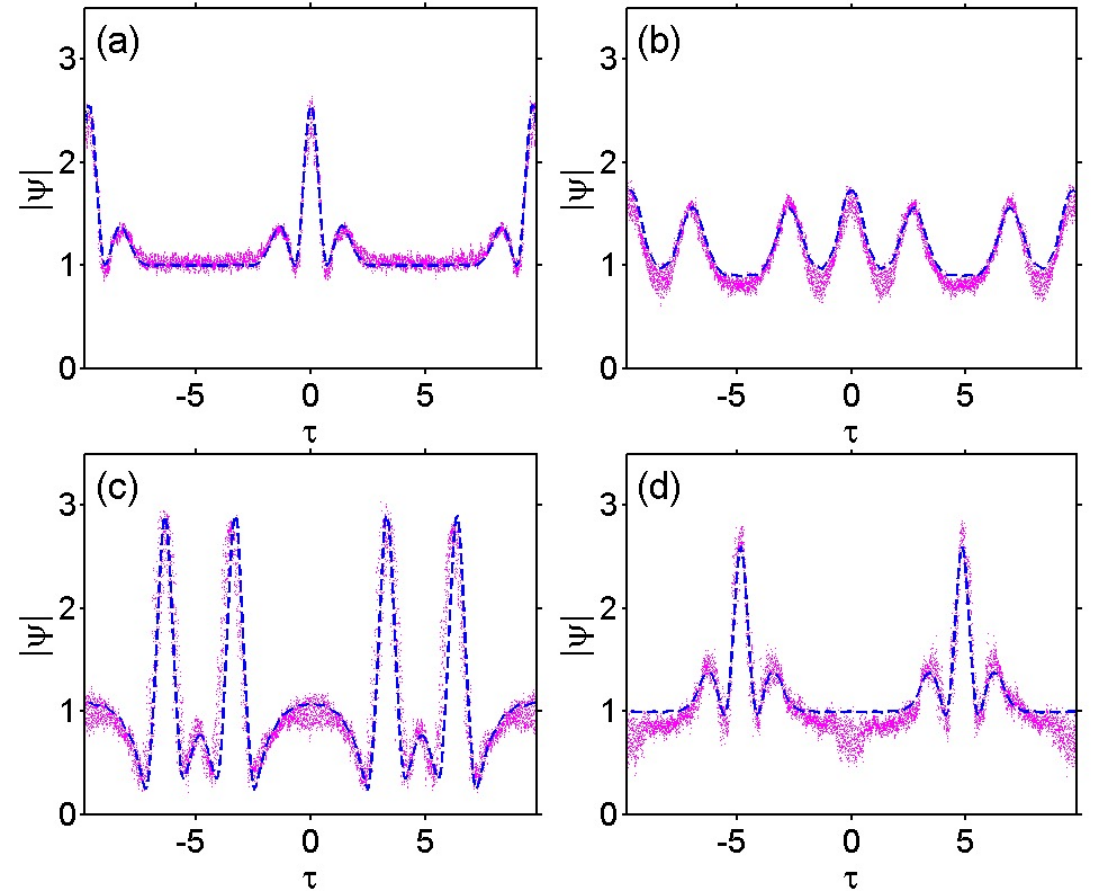
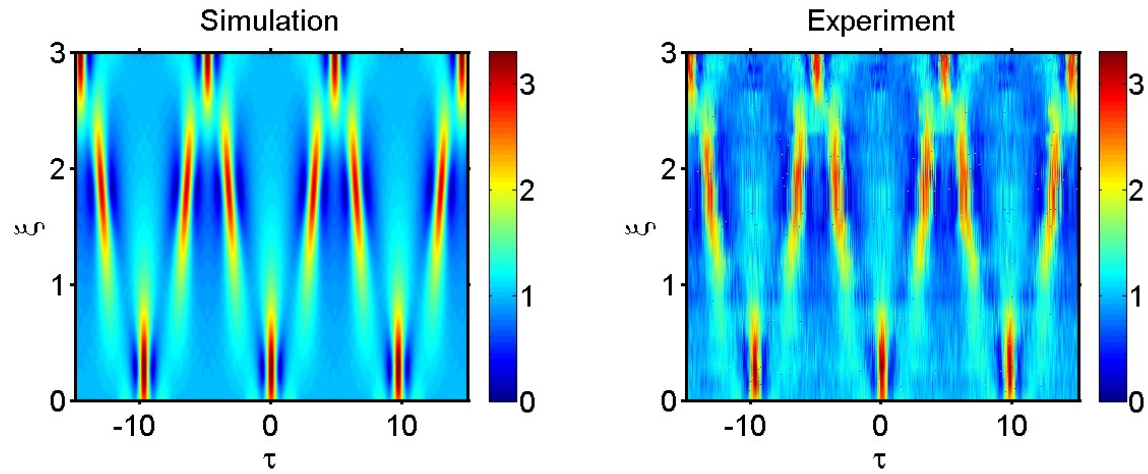
$$k = -\frac{1}{2} \left( R + \frac{1}{R} \right) \sin \alpha,$$

$$\delta x_{synch} = \frac{2\Delta x_0 + (2\Delta\theta_0 + \theta_i + \theta_f + 2\pi n) \frac{V_{gr}}{V_{ph}}}{1 - 2k \frac{V_{gr}}{V_{ph}}}$$

# Synchronization in optical fibre experiments (Dr. Bertrand Kibler, Dr. Gang Xu, university of Dijon, France) [5]



# Synchronization IIIb: recurrence



## References:

- [1] V. E. Zakharov and A.A. Gelash, Nonlinear stage of modulation instability. Phys. Rev. Lett. 111,054101, 2013.
- [2] B. Kibler, A. Chabchoub, A. Gelash, N. Akhmediev, and V. E. Zakharov. Superregular breathers in optics and hydrodynamics: Omnipresent modulation instability beyond simple periodicity. Phys. Rev. X 2015, 5, P. 041026.
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**Thank you for your attention!**

